# RIGIDITY FOR GENERAL SEMICONVEX ENTIRE SOLUTIONS TO THE SIGMA-2 EQUATION 

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#### Abstract

We show that every general semiconvex entire solution to the sigma-2 equation is a quadratic polynomial. A decade ago, this result was shown for almost convex solutions. Two decades ago, this result was obtained in three dimensions, as a byproduct of the work on special Lagrangian equations. Warren's rare saddle entire solution in 2014 confirmed the necessity of the semiconvexity assumption.


## 1. Introduction

In this paper, we show that every general semiconvex entire solution in $\mathbb{R}^{n}$ to the Hessian equation

$$
\sigma_{k}\left(D^{2} u\right)=\sigma_{k}(\lambda)=\sum_{1 \leq i_{1}<\cdots<i_{k} \leq n} \lambda_{i_{1}} \cdots \lambda_{i_{k}}=1
$$

with $k=2$ must be quadratic. Here the $\lambda_{i}$ 's are the eigenvalues of the Hessian $D^{2} u$.

## THEOREM 1.1

Let $u$ be a smooth semiconvex solution to $\sigma_{2}\left(D^{2} u\right)=1$ on $\mathbb{R}^{n}$ with $D^{2} u \geq-K I$ for a large $K>0$. Then $u$ is quadratic.

Recall the classical Liouville theorem for the Laplace equation $\sigma_{1}\left(D^{2} u\right)=$ $\Delta u=1$ or the Jörgens-Calabi-Pogorelov theorem for the Monge-Ampère equation $\sigma_{n}\left(D^{2} u\right)=\operatorname{det} D^{2} u=1$ : all convex entire solutions to those equations must be quadratic. Theorem 1.1 has been settled under an almost convexity condition $D^{2} u \geq(\delta-\sqrt{2 /[n(n-1)]}) I$ for general dimension in the joint work with Chang [4]; under the general semiconvexity condition $D^{2} u \geq-K I$ in three dimensions by taking advantage of the special Lagrangian form of the equation in this case (see [18]); or under the same general semiconvexity condition and an additional
quadratic growth assumption for general dimensions in our work [15]. Assuming a super quadratic growth condition, Bao, Chen, Guan, and Ji [1] demonstrated that all convex entire solutions to $\sigma_{k}\left(D^{2} u\right)=1$ with $k=1,2, \ldots, n$ are quadratic polynomials, and Chen and Xiang [5] showed that all "super quadratic" entire solutions to $\sigma_{2}\left(D^{2} u\right)=1$ with $\sigma_{1}\left(D^{2} u\right)>0$ and $\sigma_{3}\left(D^{2} u\right) \geq-K$ are also quadratic polynomials. Assuming only quadratic growth on entire solutions to $\sigma_{2}\left(D^{2} u\right)=1$ in three dimensions, the same rigidity result was proved in the joint work with Warren [17]. Warren's rare saddle entire solutions for the $\sigma_{2}\left(D^{2} u\right)=1$ case (see [16]) confirm the necessity of the semiconvexity or the quadratic growth assumption. It was "guessed" in the 2009 paper [4] that Theorem 1.1 should hold true.

The equation $\sigma_{2}(\kappa)=1$ prescribes the intrinsic scalar curvature of a Euclidean hypersurface $(x, u(x))$ in $\mathbb{R}^{n} \times \mathbb{R}^{1}$ with extrinsic principal curvatures $\kappa=\left(\kappa_{1}, \ldots, \kappa_{n}\right)$. The $\sigma_{2}$ function of the Schouten tensor arises in conformal geometry, and complex $\sigma_{2}$-type equations arise from the Strominger system in string theory.

Our current work, as well the previous ones [4] and [18], was inspired by Nitsche's classical paper [13], where the Legendre-Lewy transform was employed to produce an elementary proof of Jörgens' rigidity for the two-dimensional MongeAmpère equation and, in turn, Bernstein's rigidity for the two-dimensional minimal surface equation.

The Legendre-Lewy transform of a general semiconvex solution satisfies a uniformly elliptic, saddle equation. In the almost convex case (see [4]), the new equation becomes concave; thus the Evans-Krylov-Safonov theory yields the constancy of the bounded new Hessian and, in turn, the old one. To beat the saddle case, one has to be "lucky." Recall that, in general, the Evans-Krylov-Safonov theory fails, as shown by the saddle counterexamples of Nadirashvili and Vlăduţ [12]. Our earlier trace Jacobi inequality, as an alternative log-convex vehicle, other than the maximum eigenvalue Jacobi inequality, in deriving the Hessian estimates for general semiconvex solutions in [15], could rescue the saddleness. But the trace Jacobi only holds for a large enough trace of the Hessian. It turns out that the trace added by a large enough constant satisfies the elusive Jacobi inequality (Proposition 2.1)

Equivalently, the reciprocal of the shifted trace Jacobi quantity is superharmonic, and it remains so in the new vertical coordinates under the Legendre-Lewy transformation by a transformation rule (Proposition 2.2). Then the iteration arguments developed in the joint work with Caffarelli [3] show that the "vertical" solution is close to a "harmonic" quadratic at one small scale (Proposition 3.1, two steps in the execution: the superharmonic quantity concentrates to a constant in measure by applying Krylov and Safonov's weak Harnack; a variant of the superharmonic quantity, as a quotient of symmetric Hessian functions of the new potential, is very pleasantly concave and uniformly elliptic; consequently, closeness to a "harmonic" quadratic is possible by
the Evans-Krylov-Safonov theory), and the closeness improves increasingly as we rescale (this is a self-improving feature of elliptic equations; no concavity/convexity needed). Thus a Hölder estimate for the bounded Hessian is realized, and, consequently, so is the constancy of the new and then the old Hessian. See Section 3.

In closing, we remark that, in three dimensions, our proof provides a "pure" PDE way to establish the rigidity, distinct from the geometric measure theory way used in the earlier work on the rigidity for special Lagrangian equations (see [18, Theorem 1.3]).

## 2. Shifted trace Jacobi inequality and superharmonicity under Legendre-Lewy transform

Taking the gradient of both sides of the quadratic Hessian equation

$$
\begin{equation*}
F\left(D^{2} u\right)=\sigma_{2}(\lambda)=\frac{1}{2}\left[(\Delta u)^{2}-\left|D^{2} u\right|^{2}\right]=1 \tag{2.1}
\end{equation*}
$$

we have

$$
\begin{equation*}
\Delta_{F} D u=0 \tag{2.2}
\end{equation*}
$$

where the linearized operator is given by

$$
\begin{equation*}
\Delta_{F}=\sum_{i, j=1}^{n} F_{i j} \partial_{i j}=\sum_{i, j=1}^{n} \partial_{i}\left(F_{i j} \partial_{j}\right) \tag{2.3}
\end{equation*}
$$

with

$$
\begin{equation*}
\left(F_{i j}\right)=\triangle u I-D^{2} u=\sqrt{2+\left|D^{2} u\right|^{2}} I-D^{2} u>0 \tag{2.4}
\end{equation*}
$$

Here, without loss of generality, we assume that $\Delta u>0$ in what follows. Otherwise, the smooth Hessian $D^{2} u$ would be in the $\Delta u<0$ branch of the equation (2.1). Given the semiconvexity condition, the conclusion in Theorem 1.1 would be straightforward by the Evans-Krylov-Safonov theory.

The gradient square $\left|\nabla_{F} v\right|^{2}$ for any smooth function $v$ with respect to the inverse "metric" $\left(F_{i j}\right)$ is defined as

$$
\left|\nabla_{F} v\right|^{2}=\sum_{i, j=1}^{n} F_{i j} \partial_{i} v \partial_{j} v
$$

### 2.1. Shifted trace Jacobi inequality

PROPOSITION 2.1
Let $u$ be a smooth solution to $\sigma_{2}(\lambda)=1$ with $D^{2} u \geq-K I$. Set $b=\ln (\Delta u+J)$. Then we have

$$
\begin{equation*}
\Delta_{F} b \geq \varepsilon\left|\nabla_{F} b\right|^{2} \tag{2.5}
\end{equation*}
$$

for $J=8 n K / 3$ and $\varepsilon=1 / 3$.

## Proof

Step 1. Differentiation of the trace
We derive the following formulas for function $b=\ln \left(\sigma_{1}+J\right)=\ln (\Delta u+J)$ :

$$
\begin{equation*}
\left|\nabla_{F} b\right|^{2}=\sum_{i} f_{i} \frac{\left(\Delta u_{i}\right)^{2}}{\left(\sigma_{1}+J\right)^{2}} \tag{2.6}
\end{equation*}
$$

and

$$
\begin{align*}
& \Delta_{F} b \\
&= \frac{1}{\left(\sigma_{1}+J\right)}\left\{6 \sum_{i>j>k} u_{i j k}^{2}\right. \\
&\left.+\left[3 \sum_{i \neq j} u_{j j i}^{2}+\sum_{i} u_{i i i}^{2}-\sum_{i}\left(1+\frac{f_{i}}{\sigma_{1}+J}\right)\left(\Delta u_{i}\right)^{2}\right]\right\} \tag{2.7}
\end{align*}
$$

at $x=p$, where, without loss of generality, $D^{2} u(p)$ is assumed to be diagonalized and $f(\lambda)=\sigma_{2}(\lambda)$.

Noticing (2.4), it is straightforward to have the identity (2.6) and at $p$,

$$
\begin{equation*}
\Delta_{F} b=\sum_{i=1}^{n} f_{i}\left[\frac{\partial_{i i} \Delta u}{\left(\sigma_{1}+J\right)}-\frac{\left(\partial_{i} \Delta u\right)^{2}}{\left(\sigma_{1}+J\right)^{2}}\right] \tag{2.8}
\end{equation*}
$$

Next we substitute the fourth-order derivative terms $\partial_{i i} \Delta u=\sum_{k=1}^{n} \partial_{i i} u_{k k}$ in the above by lower-order derivative terms. Differentiating equation (2.2) $\sum_{i, j=1}^{n} F_{i j} \partial_{i j} u_{k}=0$ and using (2.4), we obtain at $p$,

$$
\begin{aligned}
\sum_{i=1}^{n} f_{i} \partial_{i i} \Delta u & =\Delta_{F} u_{k k}=\sum_{i, j=1}^{n} F_{i j} \partial_{i j} u_{k k}=\sum_{i, j=1}^{n}-\partial_{k} F_{i j} \partial_{i j} u_{k} \\
& =\sum_{i, j=1}^{n}-\left(\Delta u_{k} \delta_{i j}-u_{k i j}\right) u_{k i j}=\sum_{i, j=1}^{n}\left[u_{i j k}^{2}-\left(\Delta u_{k}\right)^{2}\right]
\end{aligned}
$$

Plugging the above identity in (2.8), we have at $p$,

$$
\Delta_{F} b=\frac{1}{\left(\sigma_{1}+J\right)}\left[\sum_{i, j, k=1}^{n} u_{i j k}^{2}-\sum_{k=1}^{n}\left(\Delta u_{k}\right)^{2}-\sum_{i=1}^{n} \frac{f_{i}\left(\Delta u_{i}\right)^{2}}{\sigma_{1}+J}\right] .
$$

Regrouping those terms $u_{\bigcirc \cap \mathcal{A}}, u_{\triangle \cap \bigcirc}, u_{\bigcirc \bigcirc \mathcal{O}}$, and $\Delta u_{\bigcirc}$ in the last three expressions, we obtain (2.7).

Subtracting (2.6) $* \varepsilon$ from (2.7), we have

$$
\left(\Delta_{F} b-\varepsilon\left|\nabla_{F} b\right|^{2}\right)\left(\sigma_{1}+J\right) \geq 3 \sum_{i \neq j} u_{j j i}^{2}+\sum_{i} u_{i i i}^{2}-\sum_{i}\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right)\left(\Delta u_{i}\right)^{2}
$$

with $\delta=1+\varepsilon$.
Fix $i$ and denote $t=\left(u_{11 i}, \ldots, u_{n n i}\right)$ and $e_{i}$ the $i^{\prime}$ th basis vector in $\mathbb{R}^{n}$, then the $i^{\prime}$ th term above can be written as

$$
\begin{equation*}
Q=3|t|^{2}-2\left\langle e_{i}, t\right\rangle^{2}-\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right)\langle(1, \ldots, 1), t\rangle^{2} \tag{2.9}
\end{equation*}
$$

## Step 2. Tangential projection

Equation (2.2) at $p$ yields that $t$ is tangential to the level set of the equation $\sigma_{2}(\lambda)=1$, $\langle D f, t\rangle=0$. Then by projecting $e_{i}$ and $(1, \ldots, 1)$ to the tangential space,

$$
\begin{aligned}
& E=\left(e_{i}\right)_{T}=e_{i}-\frac{f_{i}}{|D f|^{2}} D f \quad \text { and } \\
& L=(1, \ldots, 1)_{T}=(1, \ldots, 1)-\frac{(n-1) \sigma_{1}}{|D f|^{2}} D f .
\end{aligned}
$$

The coefficients of the two negative terms in the quadratic form (2.9)

$$
Q=3|t|^{2}-2\langle E, t\rangle^{2}-\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right)\langle L, t\rangle^{2}
$$

decrease, as simple symmetric computation shows that

$$
\begin{align*}
& |E|^{2}=1-\frac{f_{i}^{2}}{|D f|^{2}}<1 \\
& |L|^{2}=1-\frac{2(n-1)}{|D f|^{2}}<1, \quad \text { and }  \tag{2.10}\\
& E \cdot L=1-\frac{(n-1) \sigma_{1} f_{i}}{|D f|^{2}}
\end{align*}
$$

Step 3. Two anisotropic and nonorthogonal directions
We proceed to show that the quadratic form $Q$ is positive definite. When $t$ is perpendicular to both $E$ and $L, Q=3|t|^{2} \geq 0$. So we only need to deal with the anisotropic case, when $t$ is along $\{E, L\}$-space. The corresponding matrix of the quadratic form $Q$ is

$$
Q=3 I-2 E \otimes E-\eta L \otimes L
$$

with $\eta=1+\delta \frac{f_{i}}{\sigma_{1}+J}=1+(1+\varepsilon) \frac{f_{i}}{\sigma_{1}+J}$. The real $\xi$-eigenvector equation for (symmetric) $Q$ under nonorthogonal basis $\{E, L\}$ is

$$
\left(\begin{array}{cc}
3-2|E|^{2} & -2 E \cdot L \\
-\eta L \cdot E & 3-\eta|L|^{2}
\end{array}\right)\binom{\alpha}{\beta}=\xi\binom{\alpha}{\beta},
$$

where corresponding real eigenvalues

$$
\begin{aligned}
\xi & =\frac{1}{2}\left(\operatorname{tr} \pm \sqrt{\operatorname{tr}^{2}-4 \operatorname{det}}\right) \quad \text { with } \\
\operatorname{tr} & =6-2|E|^{2}-\eta|L|^{2} \quad \text { and } \\
\operatorname{det} & =9-6|E|^{2}-3 \eta|L|^{2}+2 \eta\left[|E|^{2}|L|^{2}-(E \cdot L)^{2}\right] .
\end{aligned}
$$

Now, by (2.10),

$$
\begin{align*}
\operatorname{tr} & =6-2\left(1-\frac{f_{i}^{2}}{|D f|^{2}}\right)-\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right)\left(1-\frac{2(n-1)}{|D f|^{2}}\right) \\
& >3-\delta \frac{f_{i}}{\sigma_{1}+J}=\frac{(3-\delta) \sigma_{1}+\delta \lambda_{i}+3 J}{\sigma_{1}+J}>0 \tag{2.11}
\end{align*}
$$

for any $\delta \leq 1.5$ and $J \geq 0$, given $\sigma_{1}=\sqrt{|\lambda|^{2}+2}>\left|\lambda_{i}\right|$ in the nontrivial remaining case.

Next, again by (2.10),

$$
\begin{aligned}
\operatorname{det}= & 6 \frac{f_{i}^{2}}{|D f|^{2}}-3 \delta \frac{f_{i}}{\sigma_{1}+J}+3\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right) \underbrace{\frac{2(n-1)}{|D f|^{2}}} \\
& +2\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right)[\frac{2(n-1) \sigma_{1} f_{i}}{|D f|^{2}}-\frac{n f_{i}^{2}}{|D f|^{2}}-\underbrace{\frac{2(n-1)}{|D f|^{2}}}] \\
> & -3 \delta \frac{f_{i}}{\sigma_{1}+J}+4\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right) \frac{(n-1) \sigma_{1} f_{i}}{|D f|^{2}} \\
& +\left[6-2 n\left(1+\delta \frac{f_{i}}{\sigma_{1}+J}\right)\right] \frac{f_{i}^{2}}{|D f|^{2}} .
\end{aligned}
$$

Then, for $\delta=1+\varepsilon=4 / 3$, we have

$$
\begin{aligned}
\operatorname{det} \cdot & \frac{\left(\sigma_{1}+J\right)|D f|^{2}}{f_{i}} \\
& \geq-3 \delta \underbrace{\left[(n-1) \sigma_{1}^{2}-2\right]}_{|D f|^{2}}+\left\{\begin{array}{c}
4\left(\sigma_{1}+J+\delta f_{i}\right)(n-1) \sigma_{1} \\
+\left[(6-2 n)\left(\sigma_{1}+J\right)-2 n \delta f_{i}\right] f_{i}
\end{array}\right\}
\end{aligned}
$$

$$
\begin{align*}
& =\left\{\begin{array}{c}
6 \delta+4(n-1) J \sigma_{1}+2(3-n) J \underbrace{f_{i}}_{\sigma_{1}-\lambda_{i}} \\
+(n-1)(4-3 \delta) \sigma_{1}^{2}+[2 n(2 \delta-1)+6-4 \delta] \sigma_{1} f_{i}-2 n \delta f_{i} \underbrace{f_{i}}_{\sigma_{1}-\lambda_{i}}
\end{array}\right\} \\
& \stackrel{\delta=4 / 3}{=} 8+2(n+1) J \sigma_{1}+2(n-3) J \lambda_{i}+\frac{2(n+1)}{3} \sigma_{1} f_{i}+\frac{8}{3} n \lambda_{i} f_{i} \tag{2.12}
\end{align*}
$$

Case $\lambda_{i} \geq 0$ : (2.13) is positive by the ellipticity $f_{i}>0$ from (2.4).
Case $0>\lambda_{i} \geq-K$ :

$$
(2.13)=2 n J \underbrace{\left(\sigma_{1}+\lambda_{i}\right)}_{\sqrt{2+|\lambda|^{2}}+\lambda_{i}>0}-6 J \lambda_{i}+2 J \sigma_{1}+\frac{8}{3} n \lambda_{i} \underbrace{f_{i}}_{\sigma_{1}-\lambda_{i}<2 \sigma_{1}}>0
$$

if $J=8 n K / 3$.
Therefore, the quadratic form $Q$ is positive definite, and we have derived the shifted Jacobi inequality (2.5) in the semiconvex case.

## Remark

In three dimensions, the Jacobi inequality (2.5) still holds for any $J \geq 0$ and $\varepsilon=1 / 3$ without the semiconvexity assumption $D^{2} u \geq-K I$. Actually, we only need to show that, in Step 3, (2.12) with $\delta=1+\varepsilon=4 / 3<1.5$ is also positive for negative $\lambda_{i}$. We would have the desired lower bound for (2.12),

$$
\operatorname{det} \cdot \frac{\left(\sigma_{1}+J\right)|D f|^{2}}{f_{i}}>8 f_{i}\left(\frac{\sigma_{1}}{3}+\lambda_{i}\right)>0
$$

if we know that $\lambda_{i}>-\sigma_{1} / 3$. Without loss of generality, we assume that $\lambda_{1} \geq \lambda_{2} \geq$ $\lambda_{3}$. Because $\lambda_{2}+\lambda_{3}=f_{1}>0$, only the smallest eigenvalue $\lambda_{3}$ could be negative. In such a negative case $\lambda_{3}=\frac{1-\lambda_{1} \lambda_{2}}{\lambda_{1}+\lambda_{2}}$ with $\lambda_{1} \lambda_{2}>1$, we do have $\lambda_{i}>-\sigma_{1} / 3$ or $\frac{\sigma_{1}}{-\lambda_{3}}>3$, because

$$
\begin{aligned}
\frac{\sigma_{1}}{-\lambda_{3}} & =-1+\frac{\left(\lambda_{1}+\lambda_{2}\right)^{2}}{\lambda_{1} \lambda_{2}-1} \\
& \geq-1+\frac{4 \lambda_{1} \lambda_{2}}{\lambda_{1} \lambda_{2}-1}>3
\end{aligned}
$$

Note that, in three dimensions, the Jacobi inequality for the log-convex $b=$ $\ln \Delta u=\ln \sqrt{2+|\lambda|^{2}}$ (with $\varepsilon=1 / 100$ ) was derived by Qiu [14, Lemma 3] for solutions to (2.1) along with variable right-hand side; the Jacobi inequality with $\varepsilon=1 / 3$
for the $\log -\max b=\ln \lambda_{\max }$ (with $\varepsilon=1 / 3$ ) was derived for solutions to (2.1) in [17, Lemma 2.2].

In general dimensions, a Jacobi inequality for sufficiently large $b=\ln u_{11}$, at points where $u_{11}=\lambda_{\max }$, was obtained for solutions having $\sigma_{3}\left(D^{2} u\right)$ lower bound to (2.1) along with variable right-hand side by Guan and Qiu [9, p. 1650]; another Jacobi inequality for sufficiently large $b=\ln \lambda_{\text {max }}$ was derived for semiconvex solutions to (2.1) in [15, Proposition 2.1], as mentioned in the introduction.

### 2.2. Superharmonicity under Legendre-Lewy transform

Set $\tilde{u}(x)=u(x)+\bar{K}|x|^{2} / 2$ for our $K$-semiconvex entire solution $u$ and, say, $\bar{K}=$ $J / n>K+1$, where $J=8 n K / 3$ is from Proposition 2.1. The $\bar{K}$-convexity of $\tilde{u}$ ensures that the smallest canonical angle of the "Lewy-sheared" "gradient" graph is larger than $\pi / 4$. This means that we can make a well-defined Legendre reflection about the origin,

$$
\begin{equation*}
(x, D \tilde{u}(x))=(D w(y), y) \in \mathbb{R}^{n} \times \mathbb{R}^{n} \tag{2.14}
\end{equation*}
$$

where $w(y)$ is the Legendre transform of $u+\frac{\bar{K}}{2}|x|^{2}$ (see [10]). Note that $y(x)=$ $D u(x)+K x$ is a diffeomorphism from $\mathbb{R}^{n}$ to $\mathbb{R}^{n}$ and

$$
0<D^{2} w=\left(D^{2} u+\bar{K} I\right)^{-1}<I
$$

More precisely, by [4, p. 663] or [15, (2.11)], the eigenvalues $\mu_{1} \leq \mu_{2} \leq \cdots \leq \mu_{n}$ of $D^{2} w$ satisfy

$$
\begin{equation*}
0<\mu_{1} \leq c(n)<1 \quad \text { and } \quad 0<c(n, K) \leq \mu_{i}<1 \quad \text { for } i \geq 2 \tag{2.15}
\end{equation*}
$$

As shown in [4, p. 663] or [15, proof of Proposition 2.4], the equation solved by the vertical coordinate Lagrangian potential $w(y)$,

$$
G\left(D^{2} w\right)=-F\left(D^{2} u\right)=-\sigma_{2}\left(-\bar{K} I+\left(D^{2} w\right)^{-1}\right)=-1,
$$

is conformally, uniformly elliptic for $K$-convex solutions $u$, in the sense that for $H_{i j}:=\sigma_{n}\left(\mu\left(D^{2} w\right)\right) G_{i j}$, the linearized operator $H_{i j} \partial_{i j}$ of equation

$$
\begin{align*}
0 & =H\left(D^{2} w\right) \\
& =\sigma_{n}\left(D^{2} w\right)\left[G\left(D^{2} w\right)+1\right] \\
& =-\sigma_{n-2}(\mu)+\underbrace{(n-1) \bar{K}}_{A_{1}} \sigma_{n-1}(\mu)-\underbrace{\left[\frac{n(n-1)}{2} \bar{K}^{2}-1\right]}_{A_{2}} \sigma_{n}(\mu) \tag{2.16}
\end{align*}
$$

is uniformly elliptic:

$$
c(n, K) I \leq\left(H_{i j}\right)=\sigma_{n}(\mu)\left(G_{i j}\right) \leq C(n, K) I .
$$

## PROPOSITION 2.2

Let $u(x)$ be a smooth solution to $\sigma_{2}(\lambda)=1$ with $D^{2} u \geq-K I$. Set

$$
a(y)=\left(\frac{1}{\mu_{1}}+\cdots+\frac{1}{\mu_{n}}\right)^{-1 / 3}=\left[\frac{\sigma_{n}}{\sigma_{n-1}}(\mu)\right]^{1 / 3}
$$

with the $\mu_{i}$ 's being the eigenvalues of the Hessian $D^{2} w(y)$ of the Legendre-Lewy transform of $u(x)+\bar{K}|x|^{2} / 2$. Then we have

$$
\triangle_{H} a \leq 0 .
$$

Proof
The trace Jacobi inequality (2.5) in Proposition 2.1 with $J=n \bar{K}$ is equivalent to

$$
\Delta_{F}(\Delta u+J)^{-1 / 3}=\Delta_{F} e^{-b / 3} \leq 0
$$

Noticing that $\Delta u+n \bar{K}=\frac{1}{\mu_{1}}+\cdots+\frac{1}{\mu_{n}}$ and applying the transformation rule from [15, Proposition 2.3], we immediately obtain the desired superharmonicity

$$
\Delta_{H} a=\sigma_{n}(\mu) \triangle_{G} a \leq 0 .
$$

As another preparation for the arguments in next section, we extend the operator $H$ in (2.16) to outside the box in (2.15), to a uniformly elliptic smooth operator with bounded $C^{1,1}$ norm $\left\|\nabla^{2} H\right\|_{L^{\infty}}$, still denoted by $H$.

## 3. Hölder Hessian estimate for saddle equation and rigidity

The Hessian bound $0<D^{2} w(y) \leq I$ ensures that establishing a local $C^{2, \alpha}$ estimate for such solutions to (2.16) will prove, by scaling, that $w(y)$ is a quadratic polynomial. By the iteration arguments developed in [3] for such smooth PDEs $H\left(D^{2} w\right)=0$ with solutions satisfying Hessian bounds, proving $C^{2, \alpha}$ regularity at a point, say the origin, reduces to showing that $w(y)$ is close to a uniform quadratic polynomial. Namely, we have the following.

## PROPOSITION 3.1

Let $u(x)$ be a smooth solution to $\sigma_{2}(\lambda)=1$ with $D^{2} u \geq-K I$ in $\mathbb{R}^{n}$. Let $w(y)$ be its Legendre-Lewy transform defined in (2.14) solving (2.16) in $\mathbb{R}^{n}$ with $0<D^{2} w \leq I$. Given any $\epsilon>0$, there exists small $\eta=\eta(n, K, \epsilon)>0$ and a quadratic polynomial $P(y)$ whose coefficients only depend on $n, K, \epsilon$, such that $H\left(D^{2} P\right)=0$ and

$$
\left|\frac{1}{\eta^{2}} w(\eta y)-P(y)\right| \leq \epsilon
$$

is valid for $|y| \leq 1$.

In the case that the level set $\left\{H\left(D^{2} w\right)=0\right\}$ were convex (in fact saddle from [4, p. 661]), the alternative way in [3] other than the Evans-Krylov-Safonov theory is the following. The Laplacian $\Delta w(y)$ is a sub- or supersolution of the linearized operator $\Delta_{H}=H_{i j} \partial^{2} / \partial y_{i} \partial y_{j}$ of $H\left(D^{2} w\right)$. The weak Harnack inequality shows that $\Delta w(y)$ concentrates in measure at a level $c$ on a small ball $B=B_{r}(0)$. Solving the equation $\Delta v=c$ on $B$ with $v=w$ on $\partial B$ furnishes the desired smooth approximation, which is uniform by the ABP estimate. The Laplacian can be replaced with any elliptic slice of the Hessian, such that the elliptic slice is a supersolution of $\Delta_{H}$, and the corresponding elliptic equation both has $C^{2, \alpha}$ interior regular solutions and allows for the ABP estimate.

However, it is not clear if the saddle level set $\left\{H\left(D^{2} w\right)=0\right\}$ of (2.16) is any of trace-convex (see [3]), max-min (see [2]), or twisted (see [6], [7]), so it is not clear if there are good PDEs which super-solve $\Delta_{H}$. Now that the remarkable superharmonic quantity $\sigma_{n}(\mu) / \sigma_{n-1}(\mu)$ in Proposition 2.2 is available, the core method in [3, pp. 687-690] becomes more realistic.

There is still one more hurdle to overcome. The superharmonic, "one-step" Hessian quotient $a^{3}=\sigma_{n}(\mu) / \sigma_{n-1}(\mu)$ is well known to be concave, but it is not uniformly elliptic, because $\sigma_{n}(\mu)$ could be arbitrarily close to zero. This prevents applying the Evans-Krylov-Safonov theory. We resolve this by substituting the concentration of $a$ into the "conformal" equation (2.16). This implies concentration of a better quantity. Observe that equation (2.16) can be written as

$$
\begin{equation*}
q(\mu):=\frac{\sigma_{n-1}(\mu)}{\sigma_{n-2}(\mu)}=\left[A_{1}-A_{2} \frac{\sigma_{n}(\mu)}{\sigma_{n-1}(\mu)}\right]^{-1} . \tag{3.1}
\end{equation*}
$$

Thus, the concentration of the higher quotient $a^{3}=\sigma_{n} / \sigma_{n-1}$ implies the concentration of the lower quotient $\sigma_{n-1} / \sigma_{n-2}$, which is also a concave operator (see [11, Theorem 15.18]). The almost-convex case, $D^{2} u \geq(-K+\delta) I$ for $K^{-2}=n(n-1) / 2$ and a $\delta>0$ considered in [4], corresponds to $A_{2}=0$. There, it was shown that (2.16) is uniformly elliptic for arbitrarily large $K$, in particular, the lower quotient $\sigma_{n-1} / \sigma_{n-2}=A_{1}^{-1}$ for $K^{-2}=n(n-1) / 2$. For arbitrary $K$, using the bound for $\mu$ in (2.15) and the result in [10, Theorem 15.18], we deduce the uniform ellipticity of $q(\mu)$ :

$$
\begin{equation*}
\partial_{\mu_{i}} q \in \frac{\sigma_{n-1}(\mu)}{\sigma_{n-2}^{3}(\mu)} \sigma_{n-2, i}^{2}(\mu)[c(n), 1] \subset[c(n, K), C(n, K)] . \tag{3.2}
\end{equation*}
$$

## Proof of Proposition 3.1

Step 1. Concentration of $a(y)$ in measure
Take positive small $\rho, \xi, \delta$ and large $k_{0}$, to be chosen later. We denote $a_{k}=$ $\min _{B_{1 / 2^{k}}} a$ with $B_{r}=B_{r}(0)$, and define a "bad set" $E_{k}=\left\{y \in B_{1 / 2^{k}}: a>a_{k}+\xi\right\}$. We claim that there is some $1 \leq l \leq k_{0}$ such that $\left|E_{\ell}\right| \leq \delta\left|B_{1 / 2^{\ell}}\right|$.

Otherwise, for all $1 \leq k \leq k_{0}$, we have $\left|E_{k}\right|>\delta\left|B_{1 / 2^{k}}\right|$. For each $k$, applying Krylov and Safonov's weak Harnack inequality (see [8, Theorem 9.22]) for supersolution $a(y)$ to $\triangle_{H} a \leq 0$ from Proposition 2.2, we have

$$
a(y) \geq \min _{B_{1 / 2^{k}}} a(y)+c(n, K) \xi \delta^{1 / p_{0}(n, K)} \quad \text { in } B_{2^{-1} 2^{-k}}
$$

That is, $a(y)$ increases by $\theta=c(n, K) \xi \delta^{1 / p_{0}}$ each time. After

$$
k_{0}=\text { the integer part of }\left[\frac{\max _{B_{1}} a(y)-\min _{B_{1}} a(y)}{\theta}+1\right] \leq \frac{n^{-1 / 3}}{\theta}+1
$$

steps, we obtain

$$
a(y) \geq \min _{B_{1 / 2} k_{0}} a(y)+\underset{B_{1}}{\operatorname{osc}} a(y)+\theta>\max _{B_{1}} a(y) \quad \text { in } B_{2^{-k_{0}-1}} .
$$

This is a contradiction. Thus our claim holds.
Before we move to Step 2, set $w_{l}(y)=2^{2 l} w\left(2^{-l} y\right)$.

## Remark

In the alternative case of the claim with $k_{0}$ replaced by $k_{0}-1$, for all $1 \leq k \leq k_{0}-1$, $\left|E_{k}\right|>\delta\left|B_{1 / 2^{k}}\right|$, after $k_{0}-1$ steps, we obtain

$$
\underset{B_{2}-k_{0}}{\text { osc }} a(y) \leq \theta<\xi \quad \text { or } \quad a_{k_{0}-1} \leq a(y)<a_{k_{0}-1}+\xi \quad \text { in } B_{2-k_{0}}
$$

Following the argument for Case 2 in [3, pp. 688-690], one can also reach the conclusion of Proposition 3.1.

Step 2. Approximation of $w_{l}(y)$ by quadratic polynomial
First, using uniform ellipticity (3.2), let us extend $\sigma_{n-1} / \sigma_{n-2}$ to a uniformly elliptic concave smooth operator $Q\left(D^{2} v\right)$ to outside the eigenvalue rectangle (2.15), $\mu \in[0, c(n)] \times[c(n, K), 1]^{n-1}$. Let $v(y) \in C^{\infty}\left(B_{1}\right)$ solve the concave equation $Q\left(D^{2} v\right)=\left(A_{1}-A_{2} a_{\ell}^{3}\right)^{-1}$ in $B_{1}$ with $v=w_{l}$ on $\partial B_{1}$. Then from the quotient representation (3.1) for the equation that $w_{l}$ solves, the ABP estimate (see [8, Theorem 9.1]) yields on $B_{1}$, and the claim in Step 1, we get

$$
\begin{aligned}
\left|w_{l}-v\right| & \leq C(n, K)\left\|Q\left(D^{2} w_{l}\right)-Q\left(D^{2} v\right)\right\|_{L^{n}\left(B_{1}\right)} \\
& \leq C(n, K) \delta^{1 / n}+C(n, K)\left\|\frac{a^{3}-a_{\ell}^{3}}{\left(A_{1}-A_{2} a^{3}\right)\left(A_{1}-A_{2} a_{\ell}^{3}\right)}\right\|_{L^{n}\left(E_{\ell}^{c}\right)} \\
& \leq C(n, K)\left(\delta^{1 / n}+\xi\right),
\end{aligned}
$$

where, in the last inequality, we used the boundedness of $\left(A_{1}-A_{2} a\right)^{-1}$ via (3.1) and (2.15). By the Evans-Krylov-Safonov theory applied to the smooth equation $Q\left(D^{2} v\right)=\left(A_{1}-A_{2} a_{\ell}^{3}\right)^{-1}, v(y)$ has uniform interior $C^{3}$ estimates, so $v$ can be replaced by its quadratic part $\bar{P}$ at the origin, up to a uniform $O\left(|y|^{3}\right)$ term. Then

$$
\left|w_{l}(y)-\bar{P}(y)\right| \leq C(n, K)\left(\delta^{1 / n}+\xi+|y|^{3}\right) .
$$

Let $y=\rho z, \overline{\bar{P}}(z)=\rho^{-2} \bar{P}(\rho z)$. For $|z| \leq 1$, we have

$$
\left|\rho^{-2} w_{l}(\rho z)-\overline{\bar{P}}(z)\right| \leq C(n, K)\left(\frac{\delta^{1 / n}+\xi}{\rho^{2}}+\rho\right)
$$

Still $H\left(D^{2} w_{l}(\rho z)\right)=0$. By [3, Lemma 2], which was proved by applying the maximum principle in a contradiction argument, and noticing that $F\left(D^{2} \overline{\bar{P}}+s I\right)=$ $H\left(D^{2} \overline{\bar{P}}+s I\right)$ is uniformly increasing as $s$ crosses $s=0$, we perturb $\overline{\bar{P}}(z)$ to another quadratic polynomial $P(z)$ so that $H\left(D^{2} P\right)=0$ with

$$
\left|\rho^{-2} w_{l}(\rho z)-P(z)\right| \leq C(n, K)\left(\frac{\delta^{1 / n}+\xi}{\rho^{2}}+\rho\right)
$$

Finally, we choose $\rho$, then $\xi, \delta$, and fix $k_{0}$ successively, depending on $n, K, \epsilon$ so that

$$
\left|\eta^{-2} w(\eta y)-P(y)\right| \leq \epsilon
$$

where $\eta=\eta(n, K, \epsilon)=\rho / 2^{l}$.

## Proof of Theorem 1.1

As indicated in the beginning of Section 2, we only need to handle the positive branch $\Delta u>0$ of the quadratic equation $\sigma_{2}\left(D^{2} u\right)=1$. This is because the only other possibility is that $D^{2} u$ is on the negative branch $\Delta u<0$ of the still elliptic and concave equation $\sigma_{2}\left(D^{2} u\right)=1$. Then the semiconvex solutions must have bounded Hessian, and consequently, the conclusion in Theorem 1.1 is straightforward by the Evans-Krylov-Safonov theory.

Now armed with Proposition 3.1, the initial closeness of $w$ to a "harmonic" quadratic on the unit ball, and repeating the proof of Proposition 2 in [3] with the equation there replaced by our smooth uniformly elliptic extended equation $H$, with bounded $C^{1,1}$ norm $\left\|\nabla^{2} H\right\|_{L^{\infty}}$, from (2.16), in the end of Section 2, we see that the closeness to "harmonic" quadratics accelerates. As in [3, p. 692], we obtain that $D^{2} w$ is Hölder at the origin. Similarly, one proves that $D^{2} w$ is Hölder in the half-ball

$$
\left[D^{2} w\right]_{C^{\alpha}\left(B_{1 / 2}\right)} \leq C(n, K)
$$

where $\alpha=\alpha(n, K)>0$.

By quadratic scaling $R^{2} w(y / R)$, we get

$$
\left[D^{2} w\right]_{C^{\alpha}\left(B_{R / 2}\right)} \leq \frac{C(n, K)}{R^{\alpha}} \longrightarrow 0 \quad \text { as } R \rightarrow \infty
$$

We conclude that $D^{2} w$ is a constant matrix, and in turn, so is $D^{2} u$.
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