

THE INVERSE PROBLEM FOR THE LOCAL GEODESIC RAY TRANSFORM

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ABSTRACT. Under a convexity assumption on the boundary we solve a local inverse problem, namely we show that the geodesic X-ray transform can be inverted locally in a stable manner; one even has a reconstruction formula. We also show that under an assumption on the existence of a global foliation by strictly convex hypersurfaces the geodesic X-ray transform is globally injective. In addition we prove stability estimates and propose a layer stripping type algorithm for reconstruction.

1. INTRODUCTION

Let X be a strictly convex domain in a Riemannian manifold (\tilde{X}, g) of dimension ≥ 3 . In this paper we consider the local inverse problem for the geodesic X-ray transform. That is, for an open set $O \subset \bar{X}$, we call geodesic segments γ of g which are contained in O with endpoints at ∂X *O-local geodesics*; we denote the set of these by \mathcal{M}_O . Thus, \mathcal{M}_O is an open subset of the smooth manifold of all geodesics, \mathcal{M} . We then define the *local geodesic transform* of a function f defined on X as the collection $(If)(\gamma)$ of integrals of f along geodesics $\gamma \in \mathcal{M}_O$, i.e. as the restriction of the X-ray transform to \mathcal{M}_O .

In order to state our main theorem in concrete terms, it is useful to introduce some notation. Let $\rho \in C^\infty(\tilde{X})$ be a defining function of ∂X , considered a function on \tilde{X} (so $\rho > 0$ in X , < 0 on $\tilde{X} \setminus \bar{X}$, vanishes non-degenerately at ∂X). Our main theorem is an invertibility result for the local geodesic transform on neighborhoods of p in \bar{X} of the form $\{\tilde{x} > -c\}$, $c > 0$, where \tilde{x} is a function with $\tilde{x}(p) = 0$, $d\tilde{x}(p) = -d\rho(p)$, see Figure 1 below.

Theorem. *For each $p \in \partial X$, there exists a function $\tilde{x} \in C^\infty(\tilde{X})$ vanishing at p and with $d\tilde{x}(p) = -d\rho(p)$ such that for $c > 0$ sufficiently small, and with $O_p = \{\tilde{x} > -c\} \cap \bar{X}$, the local geodesic transform is injective on $H^s(O_p)$, $s \geq 0$.*

Further, let $H^s(\mathcal{M}_{O_p})$ denote the restriction of elements of $H^s(\mathcal{M})$ to \mathcal{M}_{O_p} , and for $F > 0$ let

$$H_F^s(O_p) = e^{F/(\tilde{x}+c)} H^s = \{f \in H_{\text{loc}}^s(O_p) : e^{-F/(\tilde{x}+c)} f \in H^s(O_p)\}.$$

Then for $s \geq 0$ there exists $C > 0$ such that for all $f \in H_F^s(O_p)$,

$$\|f\|_{H_F^{s-1}(O_p)} \leq C \|If\|_{\mathcal{M}_{O_p}} \|H^s(\mathcal{M}_{O_p}).$$

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Remark. Here the constant C is uniform in c for small c , and indeed if we consider the regions $\{\rho \geq \rho_0\} \cap \{\tilde{x} > -c\}$ with $|\rho_0|$ and $|c|$ sufficiently small and such that this intersection is non-empty, the estimate is uniform in both c and ρ_0 .

Further, the estimate is also stable under sufficiently small perturbations of the metric g , i.e. the constant is uniform. (Notice that the hypotheses of the theorem are satisfied for small perturbations of g !)

We remark that for this result we only need to assume convexity near the point p . This local result is new even in the case that the metric is conformal to the Euclidean metric. We also point out that we also get a reconstruction method in the form of a Neumann series. See Section 2 for more details.

While this large weight $e^{F/(\tilde{x}+c)}$ means that the control over f in terms of If is weak at $\tilde{x} = -c$, the control is uniform in compact subsets of O_p : these weights are bounded below on O_p by a positive constant, and bounded above on compact subsets of O_p (in particular at parts of ∂X). Here $F > 0$ can be taken small, but not vanishing. Further, \tilde{x} , whose existence is guaranteed by the theorem, is such that $\tilde{x} = -c$ is concave from the side of O_p .

As an application, we consider domains with compact closure \overline{X} equipped with a function $\rho : \overline{X} \rightarrow [0, \infty)$ whose level sets $\Sigma_t = \rho^{-1}(t)$, $t < T$, are strictly convex (viewed from $\rho^{-1}((t, \infty))$) (and $d\rho$ is non-zero on these level sets), with $\Sigma_0 = \partial X$ and $X \setminus \cup_{t \in [0, T)} \Sigma_t = \rho^{-1}([T, \infty))$ either having 0 measure or having empty interior. (Note in particular that ρ is a boundary defining function.)

Corollary. *For X and ρ as above, if $X \setminus \cup_{t \in [0, T)} \Sigma_t$ has 0 measure, the global geodesic transform is injective on $L^2(X)$, while if it has empty interior, the global geodesic transform is injective on $H^s(X)$, $s > n/2$.*

This corollary is an immediate consequence of our main theorem. Indeed, if $If = 0$ and $f \in H^s$, $s > n/2$, $f \neq 0$, then $\text{supp } f$ has non-empty interior since f is continuous by the Sobolev embedding, while if $f \in L^2$, $f \neq 0$, then $\text{supp } f$ has non-zero measure. On the other hand, let $\tau = \inf_{\text{supp } f} \rho$; if $\tau \leq T$ we are done, for then $\text{supp } f \subset X \setminus \cup_{t \in [0, T)} \Sigma_t$. Thus, suppose $\tau > T$, so $f \equiv 0$ on Σ_t for $t < \tau$, but there exists $q \in \Sigma_\tau \cap \text{supp } f$ (since $\text{supp } f$ is closed and \overline{X} is compact). Now we use the main theorem on $\rho^{-1}(\tau, \infty)$ to conclude that a neighborhood of q is disjoint from $\text{supp } f$ to obtain a contradiction.

In fact, in this global setting we can even take $\tilde{x} = -\rho$, and the uniformity of the constants in terms of c and ρ_0 , as stated in the remark after the main theorem directly yields that if $t < T$ then there exists $\delta = \delta_t > 0$ such that if $c, \rho_0 \in (t - \delta_t, t + \delta_t)$ then a stability estimate holds (with a reconstruction method!) for the region $\rho^{-1}([\rho_0, c])$. Now in general, for $T' < T$, one can take a finite open cover of $[0, T']$ by such intervals (t'_j, t''_j) , $j = 1, \dots, k$ (with, possibly after some reindexing and dropping some intervals, $t'_1 < 0$, $t''_k > T'$, $t''_j \in (t'_{j+1}, t''_{j+1})$), and proceed inductively to recover f on $\cup_{t \in [0, T']} \Sigma_t$ from its X-ray transform, starting with the outermost region. More precisely, first, using the theorem, one can recover the restriction of f to $\rho^{-1}((-\infty, t'_1))$. Then one turns to the next interval, (t'_2, t''_2) , and notes there is a reconstruction method for the restriction to $\rho^{-1}((t'_2, t''_2))$ of functions f_2 supported in $\rho^{-1}((t'_2, +\infty))$ (no support condition needed at the other end, t''_2). One applies this to $f_2 = \phi_2 f$, where ϕ_2 identically 1 near $\rho^{-1}([t'_1, +\infty))$, supported in $\rho^{-1}((t'_2, +\infty))$; since $f = (1 - \phi_2)f + \phi_2 f$, and one has already recovered $(1 - \phi_2)f$, one also knows the X-ray transform of $\phi_2 f$, and thus the theorem is applicable.

One then proceeds inductively, covering $\rho^{-1}([0, T'])$ in k steps. This gives a *global* stability estimate, and indeed a reconstruction method doing a reconstruction layer by layer; that is, we have (in principle) developed a layer stripping algorithm for this problem.

The geodesic ray transform is closely related to the boundary rigidity problem of determining a metric on a compact Riemannian manifold from its boundary distance function. See [17, 4] for recent reviews. The case considered here is the linearization of the boundary rigidity problem in a fixed conformal class. The standard X-ray transform, where one integrates a function along straight lines, corresponds to the case of the Euclidean metric and is the basis of medical imaging techniques such as CT and PET. The case of integration along more general geodesics arises in geophysical imaging in determining the inner structure of the Earth since the speed of elastic waves generally increases with depth, thus curving the rays back to the Earth surface. It also arises in ultrasound imaging, where the Riemannian metric models the anisotropic index of refraction. Uniqueness and stability was shown by Mukhometov [11] on simple surfaces, and also for more general families of curves in two dimensions. The case of geodesics was generalized also for simple manifolds to higher dimensions in [13], [11], [1]. In dimension $n \geq 3$, the paper [3] proves injectivity and stability for the X-ray transform integrating over quite a general class of analytic curves with analytic weights, assuming an additional microlocal condition that includes the case of real-analytic metrics for a class of non-simple manifolds. Reconstruction procedures or inversion formulas have not been proven except in a few cases for instance for a class of symmetric spaces, see [5], and real-analytic curves [3]. Our results generalize support type theorems to the smooth case for the geodesic X-ray transform given in [8] for simple real-analytic metrics.

The global geometric condition that we are imposing is a natural analog of the condition $\frac{d}{dr}(r/c(r)) > 0$ proposed by Herglotz [6] and Wiechert and Zoeppritz [21] for an isotropic radial sound speed $c(r)$. In this case the geodesic spheres are strictly convex. It is also satisfied for negatively curved manifolds. But this condition allows in principle for conjugate points of the metric. In [19] one can find a microlocal study of the geodesic X-ray transform with fold caustics. A similar condition of foliating by convex hypersurfaces was used in [20] to satisfy the pseudoconvexity condition needed for Carleman estimates.

We also remark that our approach is a completely new one to uniqueness for the global problem for the geodesic ray transform. The only method up to now, except in the real-analytic category [17], has been the use of energy type equalities one introduced by Mukhometov [11] and developed by several authors which are now called ‘‘Pestov identities’’.

The main theorem is proved by considering an operator A which is essentially a ‘microlocal normal operator’ for the geodesic ray transform. Let ρ be a boundary defining function of X , i.e. $\rho > 0$ in X , $\rho = 0$ at ∂X , and $d\rho \neq 0$ at ∂X ; we assume that in fact ρ is defined on the ambient space \tilde{X} as above. First we choose an initial neighborhood U of p in \tilde{X} and a function \tilde{x} defined on it with $\tilde{x}(p) = 0$, $d\tilde{x}(p) = -d\rho(p)$, $d\tilde{x} \neq 0$ on U with convex level sets from the side of the sublevel sets and such that $O_c = \{\tilde{x} > -c\} \cap \{\rho \geq 0\}$ satisfies $\overline{O_c} \subset U$ is compact. Such a \tilde{x} exists as can be seen by slightly modifying $-\rho$, making the level sets slightly less convex. We define an operator L which integrates If over a subset of \mathcal{M}_{O_c} with a C^∞ cutoff, and consider $A = L \circ I$. We consider this operator as a map between appropriate

function spaces on O_c . It turns out that with the subset of geodesics we choose, the exponential conjugate A_F of A is a pseudodifferential operator in Melrose's scattering calculus [10]. (The exponential conjugate corresponds to working with exponentially weighted spaces for A .) We show that A_F is a Fredholm operator, and indeed that it is invertible for c near 0.

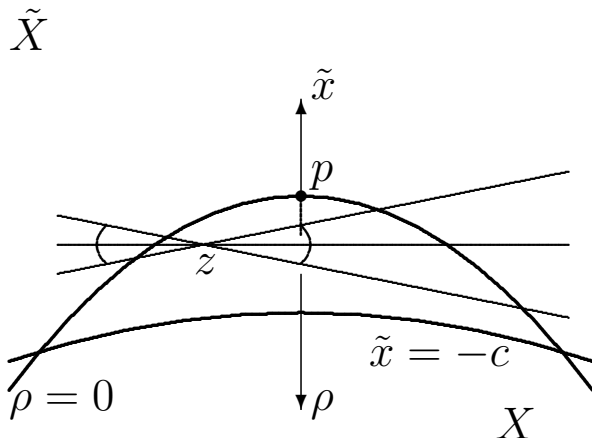


FIGURE 1. The functions ρ and \tilde{x} when the background is flat space \tilde{X} . The intersection of $\rho \geq 0$ and $x_c > 0$ (where $x_c = \tilde{x} + c$, so this is the region $\tilde{x} > -c$) is the lens shaped region O_p . Note that, as viewed from the superlevel sets, thus from O_p , \tilde{x} has concave level sets. At the point z , L integrates over geodesics in the indicated small angle. As z moves to the artificial boundary $x_c = 0$, the angle of this cone shrinks like Cx_c so that in the limit the geodesics taken into account become tangent to $x_c = 0$.

Before giving more details, recall that Stefanov and Uhlmann [18] have shown that under a microlocal condition on the geodesics, one can recover the singularities of functions from their X-ray transform, and indeed from a partial X-ray transform (where only some geodesics are included in the X-ray family \mathcal{M}'). (In fact, they also showed analogous statements for the transforms on tensors.) Roughly speaking what one needs is that given a covector $\nu = (z, \zeta)$, one needs to have a geodesic in \mathcal{M}' normal to ζ at z such that in a neighborhood of ν a simplicity condition is satisfied. Indeed, under these assumptions, a microlocal version of the normal operator, $(QI)^*(QI)$, where Q microlocalizes to \mathcal{M}' roughly speaking, is an elliptic pseudodifferential operator. Now, in dimension ≥ 3 , if the boundary ∂X is convex, one can use geodesics which are almost tangent to ∂X to give a family \mathcal{M}' which satisfies the above conditions for ν with z near ∂X . While this gives a recovery of singularities for the local problem we are considering, it yields no invertibility or reconstruction. Indeed for the latter we would like to have $(QI)^*(QI)$ to be an invertible operator on a space of functions on O_c ; in particular, as one approaches $\tilde{x} = -c$ one would need to only allow integrals over geodesics in a narrow cone, becoming tangent to $\tilde{x} = -c$, which takes one outside the framework of standard pseudodifferential operators.

To remedy this, we introduce the artificial boundary $\tilde{x} = -c$, and work with pseudodifferential operators in $x_c = \tilde{x} + c > 0$ which degenerate at $x_c = 0$. Suppressing the c dependence of x , the particular degeneration we end up with is Melrose’s scattering calculus as already mentioned. This is defined on manifolds with boundary, with boundary defining function x , and is based on degenerate vector fields $x^2\partial_x$ and $x\partial_{y_j}$, where the (x, y_1, \dots, y_{n-1}) are local coordinates. This has the effect of pushing $x = 0$ ‘to infinity’ (these vector fields are complete under the exponential map). Thus, ultimately, our approach is based on working in a framework with an artificial boundary which is effectively ‘at infinity’, and we work with function spaces allowing exponential growth at this boundary. Thus the control at $x = 0$ will be quite weak in a sense, though one has the standard control when x is bounded away from 0. Since $x = 0$ is just an artificial boundary, this is a satisfactory situation.

In fact, for most of the paper we work in a much more general setting. We consider a family of curves $\gamma_\nu : I \rightarrow \tilde{X}$ parameterized by $\nu = (z, \zeta) \in S\tilde{X}$ (the sphere bundle of \tilde{X} realized as a subbundle of $T\tilde{X}$, e.g. via a Riemannian metric) with $\gamma'_\nu(0) = \nu$ and we assume that if ν is tangent to a level set of \tilde{x} in O_c , i.e. if $\frac{d}{dt}(\tilde{x} \circ \gamma_\nu)|_{t=0} = 0$, then $\frac{d^2}{dt^2}(\tilde{x} \circ \gamma_\nu)|_{t=0} \geq C > 0$. By possibly shrinking U , we may always assume this in our setting; the lower bound on the second derivative is a concavity statement for the level sets of \tilde{x} from the side of the superlevel sets. Let $x = x_c = \tilde{x} + c$ as above. Thus, x is a boundary defining function for $\{\tilde{x} > -c\}$; for the time being we regard c as fixed. A consequence of our uniform concavity statement is that, with $\lambda = \frac{d}{dt}(x \circ \gamma_\nu)|_{t=0}$, if $C_1 > 0$ is sufficiently small and $|\lambda| < C_1\sqrt{x}$, then γ_ν remains in $x \geq 0$. Rather than using this range of λ , we instead use the stronger bound $|\lambda| < C_2x$, and define A to be an average:

$$Af(z) = x^{-1} \int If(\gamma_\nu)\chi(\lambda/x) d\mu(\nu),$$

where μ is a non-degenerate smooth measure on $S\tilde{X}$, and χ has compact support. We show that for $F \in \mathbb{R}$,

$$A_F = x^{-1}e^{-F/x}Ae^{F/x} \in \Psi_{\text{sc}}^{-1,0}(\{x \geq 0\}),$$

where Ψ_{sc} stands for the scattering calculus of Melrose, and is elliptic in the sense that the standard principal symbol is such near the boundary (up to the boundary, $x = 0$). However, even when this holds globally on a compact space, this ellipticity is not sufficient for Fredholm properties (between Sobolev spaces of order shifted by 1), or the corresponding estimates, due to the boundary $x = 0$. In general, scattering pseudodifferential operators also have a principal symbol at the boundary, which is a (typically non-homogeneous) function on a cotangent bundle; this needs to be invertible (non-zero) globally to imply Fredholm properties. Similarly, estimates implying the finite dimensionality of localized (in O) non-trivial nullspace as well as stability estimates, follow if this principal symbol is also invertible on O . (Note that here localization *does allow* the support in $\{x \geq 0\}$ to include points at $x = 0$!) We thus show that in the case when $\frac{d^2}{dt^2}(\tilde{x} \circ \gamma_\nu)|_{t=0}$ is a quadratic form in ζ subject to $\lambda = 0$, which is the case with geodesics, for suitable choices of χ , namely essentially cutoff Gaussians, this principal symbol is invertible when the weight F satisfies $F > 0$. This implies that A_F is Fredholm on this space, i.e. A itself is Fredholm on exponentially weighted spaces, where *exponential growth* is allowed at $x = 0$. We now recall that $x = x_c$ depends on c , with all estimates uniform for c remaining in

a compact set, and the argument is finished by showing that for $c > 0$ sufficiently small one not only has Fredholm properties but also invertibility, essentially as the Schwartz kernel has small support.

We note that the geodesic nature of the curves was only used in the crucial step of showing that the principal symbol at the boundary is invertible. While our argument relied on properties of the geodesics to analyze this symbol, it may well be possible to analyze it in general and prove the result for more general families of curves. We remark that J. Boman has given in [2] counterexamples for local uniqueness for the X-ray transform that integrates along lines with a dense family of smooth weights so that we expect some restrictions on the family of curves.

2. SCATTERING CALCULUS

Melrose's algebra of scattering pseudodifferential operators $\Psi_{\text{sc}}^{m,l}(\overline{M})$ on a compact manifold with boundary \overline{M} , see [10], can be thought of either via reducing to a model on \mathbb{R}^n (via appropriate charts on M , the interior of \overline{M}), or via a geometric definition. Both are of use in the current paper; the \mathbb{R}^n version makes the simplicity of this algebra transparent, while the geometric definition emphasizes that infinity in the \mathbb{R}^n -picture is not really 'remote', and indeed in our setting the artificial boundary $\tilde{x} = -c$ plays ∂M , i.e. infinity is at a decidedly finite place (moving it to infinity is what is artificial).

First we start with the \mathbb{R}^n picture, which is straightforward. Indeed, the scattering algebra in this setting is a special case of Hörmander's Weyl calculus [7, Section 18.5], which in this particular case has also been studied by Parenti [14] and Shubin [15]. That is, scattering symbols of order (m, l) are defined to be functions on $\mathbb{R}_z^n \times \mathbb{R}_\zeta^n$ satisfying

$$|D_z^\alpha D_\zeta^\beta a(z, \zeta)| \leq C_{\alpha\beta} \langle z \rangle^{l-|\alpha|} \langle \zeta \rangle^{m-|\beta|},$$

i.e. they are 'product type' symbols in z and ζ . Note that our order convention for the second order l , indicating growth/decay in z , is the opposite of that of Melrose [10] (i.e. our l is $-l$ in [10]); we make this deviation so that the symbol class increases both with m and l , i.e. so that the two indices play a parallel role. Their set is denoted by $S^{m,l}(\mathbb{R}^n, \mathbb{R}^n)$ or simply $S^{m,l}$. One then defines $\Psi_{\text{sc}}^{m,l}(\mathbb{R}^n)$ to consist of, say, left quantizations of such symbols, i.e. of operators of the form

$$(2.1) \quad Au(z) = (2\pi)^{-n} \int e^{i(z-z') \cdot \zeta} a(z, \zeta) u(z') dz' d\zeta,$$

understood as an oscillatory integral. Right quantizations could be used equally well, i.e. one gets the same class of operators if $a \in S^{m,l}$ but one substitutes $a(z', \zeta)$ into the oscillatory integral in place of $a(z, \zeta)$. Note that for $l \leq 0$, $\Psi_{\text{sc}}^{m,l}(\mathbb{R}^n)$ is a subspace of Hörmander's uniform algebra $\Psi_\infty^m(\mathbb{R}^n)$, i.e. where the above estimates hold without the factor $\langle z \rangle^{l-|\alpha|}$, and the general weight barely affects the standard arguments with pseudodifferential operators. The space $\Psi_{\text{sc}}^{*,*}(\mathbb{R}^n)$ is a filtered *-algebra under composition of operators and taking adjoints (relative to the Euclidean metric), i.e.

$$A \in \Psi_{\text{sc}}^{m,l}(\mathbb{R}^n), B \in \Psi_{\text{sc}}^{m',l'}(\mathbb{R}^n) \Rightarrow AB \in \Psi_{\text{sc}}^{m+m',l+l'}(\mathbb{R}^n)$$

and

$$A \in \Psi_{\text{sc}}^{m,l}(\mathbb{R}^n) \Rightarrow A^* \in \Psi_{\text{sc}}^{m,l}(\mathbb{R}^n).$$

Further, we define the principal symbol of A to be the equivalence class of the amplitude a in (2.1) in $S^{m,l}/S^{m-1,l-1}$, which thus captures A modulo $\Psi_{\text{sc}}^{m-1,l-1}(\mathbb{R}^n)$, i.e. one order lower operators both in terms of the differential order and growth at infinity. With this definition, the principal symbol of AB is the product of the principal symbols of A and B , while that of A^* is the complex conjugate of the principal symbol of A . In particular, if A is elliptic, i.e. its principal symbol is invertible in the sense that there is $b \in S^{-m,-l}$ such that $ab - 1 \in S^{-1,-1}$ (which is independent of the choice of representative for the principal symbol), then the standard parametrix construction produces $B \in \Psi_{\text{sc}}^{-m,-l}(\mathbb{R}^n)$ such that $AB - \text{Id} \in \Psi_{\text{sc}}^{-\infty,-\infty}(\mathbb{R}^n)$. Operators R in $\Psi_{\text{sc}}^{-\infty,-\infty}(\mathbb{R}^n)$ have a Schwartz function on \mathbb{R}^{2n} for their Schwartz kernel; this is just the inverse Fourier transform of their amplitude r in the ζ variable evaluated at $z - z'$ (where $\mathbb{R}^{2n} = \mathbb{R}_z^n \times \mathbb{R}_{z'}^n$, with z the left and z' the right variable). In particular, such operators are compact between all *polynomially weighted Sobolev spaces* $H^{s,r} = \langle z \rangle^{-r} H^s(\mathbb{R}^n)$. Further, $A \in \Psi_{\text{sc}}^{m,l}(\mathbb{R}^n)$ is bounded $H^{s,r} \rightarrow H^{s-m,r-l}$, and if A is elliptic then the parametrix construction and the compactness we observed shows that A is Fredholm – it has closed range, finite dimensional kernel and cokernel, and corresponding estimates,

$$\|u\|_{H^{s,r}} \leq C(\|Au\|_{H^{s-m,r-l}} + \|Fu\|_{H^{-N,-N}}),$$

where F can be taken a finite rank element of $\Psi_{\text{sc}}^{-\infty,-\infty}(\mathbb{R}^n)$, and N can be taken arbitrary.

In order to relate $\Psi_{\text{sc}}(\mathbb{R}^n)$ to the geometric setting, and also in order to explain its classical subalgebra, it is useful to *compactify* \mathbb{R}^n . Concretely, we compactify \mathbb{R}^n to a closed ball $\overline{\mathbb{R}^n}$ by adding the sphere at infinity \mathbb{S}^{n-1} . Thus, $\mathbb{R}^n \setminus \{0\}$ can be identified with $(0, \infty)_r \times \mathbb{S}_\theta^{n-1}$ via ‘polar coordinates’, $(r, \theta) \mapsto r\theta$; letting $x = r^{-1}$ we have ‘reciprocal polar coordinates’, $(0, \infty)_x \times \mathbb{S}_\theta^{n-1}$ which allow us to glue a sphere to $x = 0$ (corresponding to $r = \infty$) by extending the range of x to $[0, \infty)$. (Thus, formally, $\overline{\mathbb{R}^n}$ is the disjoint union of \mathbb{R}^n with $[0, \infty) \times \mathbb{S}^{n-1}$ modulo the identification of $\mathbb{R}^n \setminus \{0\}$ with $(0, \infty)_x \times \mathbb{S}_\theta^{n-1}$.) Notice that $x = r^{-1}$ is a boundary defining function near $\partial\overline{\mathbb{R}^n}$; modifying it near 0 gives a global boundary defining function ρ . It is straightforward to check that Schwartz functions on \mathbb{R}^n are exactly the restrictions to \mathbb{R}^n of C^∞ functions on $\overline{\mathbb{R}^n}$ which vanish with all derivatives at $\partial\overline{\mathbb{R}^n}$. Further, writing z as the variable on \mathbb{R}^n , the linear vector fields $z_j \partial_{z_k}$ on \mathbb{R}^n lift (automatically uniquely, as \mathbb{R}^n is the interior of $\overline{\mathbb{R}^n}$) to smooth vector fields on $\overline{\mathbb{R}^n}$ which are tangent to the boundary, and indeed all smooth vector fields tangent to the boundary are, away from the origin, linear combinations of these lifts with coefficients that are smooth on $\overline{\mathbb{R}^n}$. Since being a symbol on \mathbb{R}^n , i.e. satisfying estimates $|D_z^\alpha a(z)| \leq C_\alpha \langle z \rangle^{l-|\alpha|}$, is equivalent (away from the origin, near which one has smoothness) to satisfying stable estimates under linear vector fields, i.e. that $|V_1 \dots V_k a| \leq C \langle z \rangle^l$ for all k and linear vector fields V_j (with C depending on these), it follows that the lift of a symbol is a conormal function, i.e. a function that satisfies $\rho^l V_1 \dots V_k a \in L^\infty$ whenever V_j are vector fields tangent to $\partial\overline{\mathbb{R}^n}$, and conversely, every conormal function is the lift of a symbol. Correspondingly $\rho^{-l} C^\infty(\overline{\mathbb{R}^n}) \subset S^l(\mathbb{R}^n)$; these are the ‘classical’ or ‘one-step’ symbols; the Taylor series of a C^∞ function at the boundary gives rise to the expansion (with $x = \rho$ near $x = 0$)

$$\sum_{j \geq 0} x^{-l+j} a_j(\omega) = \sum_{j \geq 0} r^{l-j} a_j(\omega),$$

understood as an asymptotic sum.

One can now compactify *each factor* of $\mathbb{R}_z^n \times \mathbb{R}_\zeta^n$ to define the compactified space of scattering symbols $\overline{\mathbb{R}^n \times \mathbb{R}^n}$; we write ρ_∂ for the boundary defining function in the first factor ('position', z) and ρ_∞ for that in the second factor ('momentum', ζ). The same considerations as above show that a scattering symbol on $\mathbb{R}^n \times \mathbb{R}^n$ of order (m, l) corresponds to a conormal function on $\overline{\mathbb{R}^n \times \mathbb{R}^n}$, i.e. one satisfying $\rho_\infty^m \rho_\partial^l V_1 \dots V_k a \in L^\infty$ whenever V_j are vector fields tangent to both boundary hypersurfaces of $\overline{\mathbb{R}^n \times \mathbb{R}^n}$. Classical symbols, as before, then are elements of $\rho_\partial^{-l} \rho_\infty^{-m} C^\infty(\overline{\mathbb{R}^n \times \mathbb{R}^n})$, i.e. functions of the form $a = \rho_\partial^{-l} \rho_\infty^{-m} \tilde{a}$, $\tilde{a} \in C^\infty(\overline{\mathbb{R}^n \times \mathbb{R}^n})$. Note that for a classical symbol, its equivalence class in $S^{m,l}/S^{m-1,l-1}$ can be represented by $\rho_\partial^{-l} \rho_\infty^{-m}$ times the function $a_0 = \tilde{a}|_{\partial(\overline{\mathbb{R}^n \times \mathbb{R}^n})}$ on $\partial(\overline{\mathbb{R}^n \times \mathbb{R}^n})$ in the sense that any smooth extension \tilde{a}' of this function to $\overline{\mathbb{R}^n \times \mathbb{R}^n}$ produces an element of the equivalence class of a . Ellipticity then simply means the non-vanishing of this function a_0 . Note also that this principal symbol can be thought of as consisting of two parts, namely the standard principal symbol, at $\rho_\infty = 0$, and the 'boundary principal symbol' at $\rho_\partial = 0$. We also write $\Psi_{\text{sc}}^{m,l}(\overline{\mathbb{R}^n}) = \Psi_{\text{sc}}(\mathbb{R}^n)$.

If \overline{M} is a manifold with boundary with interior M , we can now define $\Psi_{\text{sc}}^{m,l}(\overline{M})$, much as the standard pseudodifferential algebra is defined on manifolds by locally identifying the manifold with \mathbb{R}^n and imposing that on such charts $U \times U$ the Schwartz kernel of the operator is that of a pseudodifferential operator on \mathbb{R}^n , and allowing additional globally smooth terms in the Schwartz kernel. In our case, the analogous construction is locally identifying \overline{M} with $\overline{\mathbb{R}^n}$, and imposing that on such charts $U \times U$ the Schwartz kernel of the operator is that of an element of $\Psi_{\text{sc}}^{m,l}(\overline{\mathbb{R}^n})$, and allowing additional globally *Schwartz* (i.e. rapidly decaying with all derivatives, smooth) terms in the Schwartz kernel. As in the standard manifold case, all the basic properties of the algebra generalize (one needs to impose some proper support conditions in the absence of compactness). Concretely, the weighted Sobolev spaces $H_{\text{sc}}^{s,r}(\overline{M})$ are also defined by local identification with $\overline{\mathbb{R}^n}$, and then $A \in \Psi_{\text{sc}}^{m,l}(\overline{\mathbb{R}^n})$ implies that A is bounded from $H_{\text{sc}}^{s,r}(\overline{M})$ to $H_{\text{sc}}^{s-m,r-l}(\overline{M})$.

It is also of some use to work out the behavior of the Schwartz kernel of elements of $\Psi_{\text{sc}}^{m,l}(\overline{M})$ on $\overline{M} \times \overline{M}$. In view of the previous definition, this reduces to a calculation for $\Psi_{\text{sc}}^{m,l}(\overline{\mathbb{R}^n})$ (modulo Schwartz terms which we ignore as they give elements of $\Psi_{\text{sc}}^{-\infty,-\infty}(\overline{M})$). Thus, using local coordinates y on \mathbb{S}^{n-1} , and corresponding coordinates (x, y, x', y') on $\overline{M} \times \overline{M}$, one checks that in the coordinates

$$x, y, X = \frac{x - x'}{x^2}, Y = \frac{y - y'}{x},$$

valid for $x > 0$, so the diagonal is $X = 0, Y = 0$ when $x > 0$, the Schwartz kernel of an element of $\Psi_{\text{sc}}^{m,l}$ is of the form $x^{-l} \tilde{K}$, where \tilde{K} is smooth in (x, y) down to $x = 0$ with values in conormal distributions on $\mathbb{R}_{X,Y}^n$, conormal to $\{X = 0, Y = 0\}$, which are Schwartz at infinity (i.e. decay rapidly at infinity with all derivatives). Further, the boundary principal symbol is simply x^{-l} times the Fourier transform in (X, Y) of $\tilde{K}|_{x=0}$ (a restriction which makes sense in view of the stated smoothness). In particular, when $l = 0$, we need to check that

$$\int e^{-i\xi X - i\eta \cdot Y} \tilde{K}(0, y, X, Y) dX dY$$

is a non-zero function of (y, ξ, η) , with a lower bound $C(\langle \xi, \eta \rangle)$, $C > 0$, for its absolute value (which means we also need a uniform bound at infinity in addition

to the invertibility). Checking this will be the main step of the arguments presented in the next section.

We mention here that vector fields in $\mathcal{V}_{\text{sc}}(\overline{M}) = x\mathcal{V}_{\text{b}}(\overline{M})$, where $\mathcal{V}_{\text{b}}(\overline{M})$ is the set of all smooth vector fields tangent to ∂M , are in $\Psi_{\text{sc}}^{1,0}(\overline{M})$, and indeed the Sobolev spaces of positive integer differential orders s are equivalently defined by $u \in H_{\text{sc}}^{s,r}(\overline{M})$ if and only if $x^{-r}V_1 \dots V_k u \in L_{\text{sc}}^2(\overline{M})$ for all $k \leq s$ (including $k = 0$) and $V_j \in \mathcal{V}_{\text{sc}}(\overline{M})$; here $L_{\text{sc}}^2(\overline{M})$ is the L^2 space given by identification by $\overline{\mathbb{R}^n}$, i.e. the measure (or density) is, up to a non-degenerate positive multiple, $r^{n-1} dr dy = x^{-n-1} dx dy$. (Densities like this may be called *scattering densities*.)

We now briefly relate the standard Sobolev spaces $H^s(\overline{M})$ to $H_{\text{sc}}^{s,r}(\overline{M})$ for $s \geq 0$. First, for $s = 0$, the above description gives $H_{\text{sc}}^{0, -(n+1)/2}(\overline{M}) = L^2(\overline{M})$ (in the sense of equivalent norms). Next, using that $V_1' \dots V_k' u \in L^2(\overline{M})$ for $k \leq s$ and V_j' smooth vector fields on \overline{M} (which is equivalent to $u \in H^s(\overline{M})$) implies that $V_1 \dots V_k u \in H_{\text{sc}}^{0, -(n+1)/2}(\overline{M})$ for $k \leq s$ and $V_j \in \mathcal{V}_{\text{sc}}(\overline{M})$ (since all elements of $\mathcal{V}_{\text{sc}}(\overline{M})$ are smooth vector fields), i.e. that $u \in H_{\text{sc}}^{s, -(n+1)/2}(\overline{M})$, so

$$(2.2) \quad H^s(\overline{M}) \subset H_{\text{sc}}^{s,r}(\overline{M}), \quad r \leq -\frac{n+1}{2},$$

with continuous inclusion map. For the converse direction, we note that if V' is a smooth vector field, then $x^2 V' \in \mathcal{V}_{\text{sc}}(\overline{M})$. Thus, $V_1 \dots V_k u \in H_{\text{sc}}^{0, 2s-(n+1)/2}(\overline{M})$ for $k \leq s$ and $V_j \in \mathcal{V}_{\text{sc}}(\overline{M})$, so $x^{-2s} V_1 \dots V_k u \in H_{\text{sc}}^{0, -(n+1)/2}(\overline{M})$, so $x^{-2} V_1 \dots x^{-2} V_k u \in H_{\text{sc}}^{0, -(n+1)/2}(\overline{M})$, implies that $V_1' \dots V_k' u \in L^2(\overline{M})$ for $k \leq s$ and V_j' smooth vector fields. Thus,

$$(2.3) \quad H_{\text{sc}}^{s,r}(\overline{M}) \subset H^s(\overline{M}), \quad r \geq -\frac{n+1}{2} + 2s,$$

with continuous inclusion map. There are similar inclusions between negative order spaces. For instance, as $H^{-s}(\overline{M}) = (H_0^s(\overline{M}))^*$, $s \geq 0$, via identification by the L^2 pairing, and as $H_0^s(\overline{M})$ is a closed subspace of $H^s(\overline{M})$, the inclusion (2.2) gives the continuous inclusion map on the dual spaces

$$(2.4) \quad H_{\text{sc}}^{-s,-r}(\overline{M}) \subset H^{-s}(\overline{M}), \quad -r \geq \frac{n+1}{2}.$$

Finally we discuss what happens when ellipticity holds only locally. Thus, suppose O is an open subset of \overline{M} on which $A \in \Psi_{\text{sc}}^{m,l}(\overline{M})$ is elliptic, and suppose that $K \subset O$ is a compact subset. Let ϕ be supported in O , identically 1 on K ; let O' be a neighborhood of $\text{supp } \phi$ with closure compactly contained in O . By the ellipticity assumption, there is a *local parametrix* $G \in \Psi_{\text{sc}}^{-m,-l}(\overline{M})$ for A such that $GA = \text{Id} + E$, $E \in \Psi_{\text{sc}}^{0,0}(\overline{M})$, but over O' the better conclusion that E is, locally, in $\Psi_{\text{sc}}^{-\infty,-\infty}$, holds, so $\phi E \phi \in \Psi_{\text{sc}}^{-\infty,-\infty}(\overline{M})$. Thus, $\phi E \phi$ is compact on any polynomially weighted Sobolev space, so in particular there is a finite rank operator $F \in \Psi_{\text{sc}}^{-\infty,-\infty}(\overline{M})$ supported in $O \times O$ such that $\text{Id} + \phi E \phi - F$ is invertible. Now suppose that v is supported in K , so $\phi v = v$. Then $\phi G A \phi = \phi^2 + \phi E \phi$ shows that

$$(\text{Id} + \phi E \phi)v = \phi G A v,$$

so

$$v = (\text{Id} + \phi E \phi - F)^{-1} \phi G A v - (\text{Id} + \phi E \phi - F)^{-1} F v.$$

In particular, if $A v = 0$ then v is in a finite dimensional space, namely the range of $(\text{Id} + \phi E \phi - F)^{-1} F$, and if one chooses a complementary subspace V of $\text{Ker } A \cap \{w :$

$\text{supp } w \subset K\}$ in a weighted Sobolev space $H_{\text{sc}}^{s,r}(\overline{M})$, then there is a constant $C > 0$ such that for $v \in V \cap \{w : \text{supp } w \subset K\}$,

$$\|v\|_{H_{\text{sc}}^{s,r}(\overline{M})} \leq C \|Av\|_{H_{\text{sc}}^{-m,r-l}(\overline{M})},$$

i.e. a stability estimate holds.

Now suppose that one has a family of operators, A_t , $t \in [0, T]$, depending continuously on t in $\Psi_{\text{sc}}^{m,l}(\overline{M})$, with each element of the family being elliptic on O (and thus there is a uniform constant in the estimates over compact subsets of O). Suppose also that we have a continuous function f on $[0, T]$ with $f(0) = 0$, a compact subset K_0 of O , and a family of open sets \mathcal{O}_t , $t > 0$ in K_0 , with the boundary defining function satisfying $x \leq f(t)$ on \mathcal{O}_t , and we are interested in distributions v supported in \mathcal{O}_t . In view of the uniform elliptic estimates, choosing O' a neighborhood of K_0 with closure compactly contained in O , we then have families of operators G_t and E_t , depending continuously on $t \in [0, T]$, with values in $\Psi_{\text{sc}}^{-m,-l}(\overline{M})$, resp. $\Psi_{\text{sc}}^{0,0}(\overline{M})$, such that on O' , E_t is uniformly in $\Psi_{\text{sc}}^{-\infty,-\infty}(\overline{M})$. Thus, the Schwartz kernel K_t of E_t satisfies that for any N , $x^{-N}(x')^{-N}K_t$ is bounded (with values in scattering densities in the right, i.e. primed, factor), i.e. locally is of the form $\kappa_t \frac{dx' dy'}{(x')^{n+1}}$ with $|\kappa_t(x, y, x', y')| \leq C_N x^N (x')^N$. (Notice that the fact that we used ‘scattering’ densities is thus of little relevance; any polynomial factor such as $(x')^{-n-1}$, can make no difference.) If $\phi_t \in C_c^\infty(\overline{M})$ is now supported in \mathcal{O}_t and takes values in $[0, 1]$, then $\phi_t E_t \phi_t$ has kernel $\phi_t(x, y) \phi_t(x', y') \kappa_t \frac{dx' dy'}{(x')^{n+1}}$, with $|\phi_t(x, y) \phi_t(x', y') \kappa_t| \leq C'_N f(t)^{2N} x^{n+1} (x')^{n+1}$ for all N , and thus by Schur’s lemma is bounded on $L_{\text{sc}}^2(\overline{M})$ with norm $\leq C''_N f(t)^{2N}$. In particular, there is $t_0 > 0$ such that the norm is $< 1/2$ for $t \in (0, t_0]$. Thus, $\text{Id} + \phi_t E_t \phi_t$ is invertible for such t , and the previous arguments give that if $K_t \subset \mathcal{O}_t$ is compact then for $t \in (0, t_0]$,

$$\text{Ker } A_t \cap \{w : \text{supp } w \subset K_t\} = \{0\}$$

and for v supported in K_t one has the stability estimate (with uniform constant C)

$$\|v\|_{H_{\text{sc}}^{s,r}(\overline{M})} \leq C \|A_t v\|_{H_{\text{sc}}^{-m,r-l}(\overline{M})}.$$

We remark here that $(\text{Id} + \phi_t E_t \phi_t)^{-1}$ can be constructed by a Neumann series, and thus ultimately our whole argument is completely constructive.

In our setting we start with an ambient manifold \tilde{X} with equipped with a function \tilde{x} with non-degenerate level sets near the 0 value, let $x_c = \tilde{x} + c$ (c near 0), let $M_c = \{x_c > 0\}$, identify a neighborhood of $Y = \{\tilde{x} = 0\}$ with $Y \times (-\delta, \delta)_{\tilde{x}}$, and have a family of operators $B_c \in \Psi_{\text{sc}}^{m,l}(\overline{M}_c)$ with Schwartz kernel localized in $\tilde{x} < c_0$ (in both factors), where $c_0 > 0$ is small. We further have a fixed set $O \subset \tilde{X}$ with compact closure, $K \subset O$ compact, and a function f continuous on $[0, \delta)$ with $f(0) = 0$, such that on $O \cap M_c$, $x_c \leq f(c)$. In order to analyze the B_c as $c \rightarrow 0$, we regard these instead as operators on $M_0 = \{\tilde{x} > 0\}$ by letting $A_c = (\Phi_c^{-1})^* B_c \Phi_c^*$, $\Phi_c(\tilde{x}, y) = (\tilde{x} + c, y)$ which maps M_c to M_0 . The operators A_c obtained by this procedure (with the parameter being c rather than t), together with the corresponding translates \mathcal{O}_c and K_c of $O \cap M_c$ and $K \cap M_c$ satisfy all the requirements of the previous paragraphs, and thus conclusions apply, which, when translated to B_c give that for sufficiently small c

$$\text{Ker } B_c \cap \{w \in H_{\text{sc}}^{s,r}(M_c) : \text{supp } w \subset K \cap M_c\} = \{0\}$$

and for $v \in H_{\text{sc}}^{s,r}(M_c)$ supported in K one has the stability estimate (with uniform constant C)

$$\|v\|_{H_{\text{sc}}^{s,r}(\overline{M_c})} \leq C \|B_c v\|_{H_{\text{sc}}^{s-m, r-l}(\overline{M_c})}.$$

Further, in our setting, the operators B_c are in fact of the form

$$B_c = x_c^{-1} e^{-F/x_c} L_c e^{F/x_c} \in \Psi_{\text{sc}}^{-1,0}(\overline{M_c}),$$

so we in fact obtain that for sufficiently small c

$$\text{Ker } L_c \cap \{e^{-F/x_c} w \in H_{\text{sc}}^{s,r}(M_c) : \text{supp } w \subset K \cap M_c\} = \{0\}$$

and for $v \in e^{F/x_c} H_{\text{sc}}^{s,r}(M_c)$ supported in K one has the stability estimate (with uniform constant C)

$$\|e^{-F/x_c} v\|_{H_{\text{sc}}^{s,r}(\overline{M_c})} \leq C \|e^{-F/x_c} L_c v\|_{H_{\text{sc}}^{s+1, r-1}(\overline{M_c})}.$$

Notice that this is an exponentially weak estimate at ∂M_c , i.e. at $x_c = 0$, but the exponential factor is immaterial in $x_c > 0$. Notice also that if $\tilde{v} \in H^s(\tilde{X})$, say, then for $F > 0$ its restriction v to M_c is in $e^{F/x_c} H_{\text{sc}}^{s,r}(M_c)$ for all r , i.e. the results are in fact applicable to v .

3. PROOFS

Suppose first that X is a domain in (\tilde{X}, g) , $p \in \partial X$, and ∂X is geodesically strictly convex at p (hence near p). That is, with ρ a boundary defining function of \tilde{X} , we have (with G the dual metric, and metric function) that for covectors $\beta \in T_p^* \tilde{X} \setminus o$,

$$(H_G \rho)(\beta) = 0 \Rightarrow (H_G^2 \rho)(\beta) < 0.$$

In particular, by compactness of the unit sphere and homogeneity, there is a neighborhood U_0 of p in \tilde{X} and $C_0 > 0$, $\delta > 0$ such that for covectors $\beta \in T_{U_0}^* \tilde{X} \setminus o$,

$$|(H_G \rho)(\beta)| < \delta G(\beta)^{1/2} \Rightarrow (H_G^2 \rho)(\beta) \leq -C_0 G(\beta).$$

We then want to define a function \tilde{x} near p such that $\tilde{x}(p) = 0$, the region $\tilde{x} \geq -c$, $\rho \geq 0$, is compact for $c > 0$ small, and the level sets of \tilde{x} are concave from the side of this region (i.e. the super-level sets of \tilde{x}). By shrinking U_0 if needed, we may assume that it is a coordinate neighborhood of p . Concretely we let, for $\epsilon > 0$ to be decided, an with $|\cdot|$ the Euclidean norm,

$$\tilde{x}(z) = -\rho(z) - \epsilon |z - p|^2;$$

then $\tilde{x} \geq -c$ gives $\rho + \epsilon |z - p|^2 \leq c$ and thus $\rho \leq c$; further, with $\rho \geq 0$ this gives $|z - p|^2 \leq c/\epsilon$. Thus, for c/ϵ sufficiently small, the region $\tilde{x} \geq -c$, $\rho \geq 0$, is compactly contained in U_0 . Further, for $\beta \in T_{U_0}^* \tilde{X}$, $H_G \tilde{x}(\beta) = -H_G \rho(\beta) - \epsilon H_G |z - p|^2$, so $H_G \tilde{x} = 0$ implies $|H_G \rho| < C' \epsilon G^{1/2}$, so with $\delta > 0$ as above there is $\epsilon' > 0$ such that for $\epsilon \in (0, \epsilon')$, $H_G \tilde{x} = 0$ in U_0 implies $|H_G \rho| < \delta G^{1/2}$, and then, for $\epsilon < \epsilon'$,

$$H_G^2 \tilde{x} = -H_G^2 \rho - \epsilon H_G^2 |z - p|^2 \geq (C_0 - C'' \epsilon) G.$$

Thus, there is $\epsilon_0 > 0$ such that for $\epsilon \in (0, \epsilon_0)$, $H_G^2 \tilde{x} \geq (C_0/2)G$ at $T_p^* \tilde{X}$ when $H_G \tilde{x}$ vanishes. Thus taking $c_0 > 0$ sufficiently small (corresponding to ϵ_0), we have constructed a function \tilde{x} defined on a neighborhood U_0 of p with concave level sets (from the side of the super-level sets) and such that for $0 \leq c \leq c_0$,

$$O_c = \{\tilde{x} > -c\} \cap \{\rho \geq 0\}$$

has compact closure in $U_0 \cap \tilde{X}$.

From now on we work with $x_c = \tilde{x} + c$, which is the boundary defining function of the region $x_c \geq 0$; we suppress the c dependence and simply write x in place of x_c . For most of the following discussion we completely ignore the actual boundary, $\rho = 0$; this will only play a role at the end since ellipticity properties only hold in U_0 and we need f to be supported in $\rho \geq 0$, ensuring localization, in order to obtain injectivity and stability estimates. Thus, completing \tilde{x} to a coordinate system (\tilde{x}, y) on a neighborhood $U_1 \subset U_0$ of p , for each point (\tilde{x}, y) we can parameterize geodesics through this point by the unit sphere; the relevant ones for us are ‘almost tangent’ to level sets of \tilde{x} , i.e. we are interested in ones with tangent vector $c(\lambda\partial_x + \omega\partial_y)$, $c > 0$ (to say have unit length), $\omega \in \mathbb{S}^{n-2}$, and λ relatively small.

Now, the geodesic corresponding to (z_0, ν_0) , $\gamma = \gamma_{z_0, \nu_0}$, is the projection of the bicharacteristic $\tilde{\gamma}$ emanating from $(z_0, g_{z_0}(\nu_0)) = (z_0, \zeta_0)$ (i.e. the integral curve of H_G through this point; here we are using the metric g_{z_0} to turn the vector ν_0 into a covector) which thus satisfies $(\frac{d}{dt}\tilde{\gamma})(t) = H_G(\tilde{\gamma}(t))$, so $\frac{d^2}{dt^2}(f \circ \tilde{\gamma})(t) = H_G^2 f(\tilde{\gamma}(t))$. Thus, if f is a function on the base space \tilde{X} then $(\frac{d^2}{dt^2}\gamma)(0) = (H_G^2 f)(\gamma(0), g_{\gamma(0)}(\gamma'(0)))$. But H_G^2 is homogeneous degree two in the fiber (second) variable of the cotangent bundle, and it is a polynomial, which shows that $\frac{d^2}{dt^2}(f \circ \tilde{\gamma})(0)$ is a quadratic polynomial in ν .

We now make this more concrete. For this, we use a fibration by level sets of a function x with non-vanishing differential. Letting V be a vector field orthogonal with respect to g to these level sets with $Vx = 1$, and using $\{x = 0\}$ as the initial hypersurface, the flow of V (locally) identifies a neighborhood of $\{x = 0\}$ with $(-\epsilon, \epsilon)_x \times \{x = 0\}$, with the first coordinate being exactly the function x (since time t flow by V changes the value of x by t). In particular, choosing coordinates y_j on $\{x = 0\}$, we obtain coordinates on this neighborhood such that ∂_{y_j} and ∂_x are orthogonal, i.e. the metric is of the form $f(x, y) dx^2 + h(x, y, dy)$, and the dual metric is of the form

$$F(x, y)\xi^2 + \sum H_{ij}(x, y)\eta_i\eta_j,$$

with $f, F > 0$, so (with h_{ij} denoting the metric components, so H_{ij} is the inverse matrix of h_{ij}),

$$\frac{dx}{dt} = 2F(x, y)\xi, \quad \frac{dy_i}{dt} = 2 \sum H_{ij}(x, y)\eta_j, \quad -\frac{d\xi}{dt} = \frac{\partial F}{\partial x}\xi^2 + \sum \frac{\partial H_{ij}}{\partial x}(x, y)\eta_i\eta_j,$$

and thus

$$\begin{aligned} \frac{1}{2} \frac{d^2 x}{dt^2} &= 2 \frac{\partial F}{\partial x}(x, y)F(x, y)\xi^2 + 2 \sum \frac{\partial F}{\partial y_i} H_{ij}(x, y)\eta_j \xi \\ &\quad - F(x, y) \frac{\partial F}{\partial x}\xi^2 - F(x, y) \sum \frac{\partial H_{ij}}{\partial x}(x, y)\eta_i\eta_j \end{aligned}$$

which at $\frac{dx}{dt} = 0$, thus $\xi = 0$, simplifies to

$$\begin{aligned} & - \sum \frac{\partial H_{ij}}{\partial x}(x, y)h_{ik}(x, y)h_{jl}(x, y) \frac{dy_k}{dt} \frac{dy_l}{dt} \\ &= - \sum \frac{\partial H_{ij}}{\partial x}(x, y)h_{ik}(x, y)h_{jl}(x, y)\omega_k\omega_l. \end{aligned}$$

Here we used the unit sphere for the ω -parameterization. Note that

$$-(\sum \frac{\partial H_{ij}}{\partial x}(x, y)h_{ik}(x, y)h_{jl}(x, y))_{kl}$$

is positive definite by our assumptions. Thus, for geodesics we have a positive definite quadratic form

$$\alpha(x, y, \omega, 0, 0) = - \sum \left(\sum \frac{\partial H_{ij}}{\partial x}(x, y) h_{ik}(x, y) h_{jl}(x, y) \right)_{kl} \omega_k \omega_l.$$

In fact, as explained in the introduction, we mostly work in the following more general setting. We consider integrals along a family of C^∞ curves $\gamma_{x,y,\lambda,\omega}$ in \mathbb{R}^n , $(x, y, \lambda, \omega) \in \mathbb{R} \times \mathbb{R}^{n-1} \times \mathbb{R} \times \mathbb{S}^{n-2}$, depending smoothly (C^∞) on the parameters, typically (but not necessarily) geodesics. Here \mathbb{R}_y^{n-1} could be replaced by an arbitrary manifold and below we make x small, so effectively we are working in a tubular neighborhood of a codimension one submanifold of an arbitrary manifold, such as \tilde{X} . However, since the changes in the manifold setting are essentially just notational, for the sake of clarity we work with \mathbb{R}^n . Further, below we work with neighborhoods of a compact subset $\{0\} \times K \subset \mathbb{R}_x \times \mathbb{R}_y^{n-1}$; $\gamma_{x,y,\lambda,\omega}(t)$ would only need to be defined for (x, y) in a fixed neighborhood \tilde{U} of $\{0\} \times K$ and for $|\lambda| < \tilde{\delta}_0$, and $|t| < \tilde{\delta}_0$, $\tilde{\delta}_0 > 0$ a fixed constant.

The basic feature we need is that for $x \geq 0$ and for λ sufficiently small, depending on x , the curves stay in $[0, \infty) \times \mathbb{R}^{n-1}$. Thus, for $x = 0$ only the parameter value $\lambda = 0$ is allowed; in our concrete setting $|\lambda| \leq C_0 \sqrt{x}$ works for suitably small $C_0 > 0$. However, it is convenient to use an even smaller range of λ , such as $|\lambda| \leq C_0 x$. So concretely assume that

$$\begin{aligned} \gamma_{x,y,\lambda,\omega}(0) &= (x, y), \quad \gamma'_{x,y,\lambda,\omega}(0) = (\lambda, \omega), \\ \gamma''_{x,y,\lambda,\omega}(t) &= 2(\alpha(x, y, \lambda, \omega, t), \beta(x, y, \lambda, \omega, t)), \end{aligned}$$

and

$$\alpha(0, y, 0, \omega, 0) \geq 2C > 0,$$

with α, β smooth. This implies that if $K \subset \mathbb{R}^{n-1}$ is compact, then for a sufficiently small neighborhood U of $\{0\} \times K$ in \mathbb{R}^n (with compact closure), and for λ and t sufficiently small, say $|\lambda|, |t| < \delta_0$, where $\delta_0 > 0$, one has

$$\alpha(x, y, \lambda, \omega, t) \geq C > 0.$$

One may assume that $x < \delta_0$ on U . Thus, writing $\gamma(t) = (x'(t), y'(t))$,

$$x' = x + \lambda t + t^2 \int_0^1 (1-s) \alpha(x, y, \lambda, \omega, s) ds \geq x + \lambda t + Ct^2/2,$$

so if $|t| < \delta_0$, $(x, y) \in U$, $|\lambda| < \delta_0$ then

$$(3.1) \quad x' \geq \frac{C}{2} \left(t + \frac{\lambda}{C} \right)^2 + \left(x - \frac{\lambda^2}{2C} \right).$$

Thus, for $|\lambda| \leq \sqrt{2C} \sqrt{x}$ (and $|\lambda| < \delta_0$), $x' \geq 0$, i.e. the curves remain in the half-space $x' \geq 0$ at least for $|t| < \delta_0$. Further, if we fix $x_0 > 0$, then $x' \geq x_0$ provided $|t + \frac{\lambda}{C}| > \sqrt{2x_0/C}$ and $|t| < \delta_0$, thus when $|\lambda| \leq C_0 x_0$ and $|\lambda| < \delta_0$ then $x' \geq x_0$ provided $|t| > \frac{C_0}{C} x + \sqrt{2x_0/C}$, $|t| < \delta_0$. Assuming $x \leq x_0$ and taking x_0 sufficiently small so that $\frac{C_0}{C} x_0 + \sqrt{2x_0/C} < \delta_0$, we thus deduce that the curve segments $\gamma_{x,y,\lambda,\omega}|_{(-\delta_0, \delta_0)}$ are outside the region $x' < x_0$ for t outside a (fixed!) compact subinterval of $(-\delta_0, \delta_0)$. From now on, by γ we mean the restriction $\gamma_{x,y,\lambda,\omega}|_{(-\delta_0, \delta_0)}$, and we everywhere assume that the functions we integrate along γ are supported in $x' \leq x_0/2$, so all integrals are on a fixed compact subinterval.

Note that in the case of geodesics, as discussed above, α is a quadratic polynomial in w ; this will be of use when the ellipticity of the boundary principal symbol is discussed.

Before we proceed, we discuss the blowup of a space around a submanifold. Here we work locally on say \mathbb{R}_w^m , thus the submanifold can be taken to be given by $w' = 0$, where we write $w = (w', w'') \in \mathbb{R}^k \times \mathbb{R}^{m-k}$. Then blowing up $\mathbb{R}^{m-k} = \{w' = 0\}$ in \mathbb{R}^m amounts to introducing cylindrical coordinates around it, i.e. the factor \mathbb{R}^{m-k} (the cylindrical ‘axis’, though higher dimensional) is unchanged, while on $\mathbb{R}_{w'}^k$, one introduces ‘polar coordinates’ $(|w'|, \frac{w'}{|w'|}) \in [0, \infty) \times \mathbb{S}^{k-1}$, thus one replaces \mathbb{R}^m by

$$[\mathbb{R}^m; \mathbb{R}^{m-k}] = [0, \infty)_r \times \mathbb{S}_\theta^{k-1} \times \mathbb{R}_{w''}^{m-k};$$

altogether one has ‘coordinates’ (the quotes are due to the spherical factor)

$$r = |w'|, \quad \theta = \frac{w'}{|w'|}, \quad w'',$$

with the equalities holding outside $r = 0$. The new boundary

$$\text{ff} = \{0\} \times \mathbb{S}^{k-1} \times \mathbb{R}^{m-k}$$

is called the front face. Further one has a blow-down map $\Phi : [\mathbb{R}^m; \mathbb{R}^{m-k}] \rightarrow \mathbb{R}^m$ which is smooth, namely $(r, \theta, w'') \mapsto (r\theta, w'')$, but is not invertible at $r = 0$ although it restricts to a diffeomorphism $[\mathbb{R}^m; \mathbb{R}^{m-k}] \setminus \text{ff} \rightarrow \mathbb{R}^m \setminus \mathbb{R}^{m-k}$. We refer to the Appendix of [10] for a concise but more detailed description, and for further references. Note that the effect of this blow up is to distinguish directions of approach to the submanifold being blown up, \mathbb{R}^{m-k} ; curves $c : [0, \epsilon) \rightarrow \mathbb{R}^m$ with $c(0) \in \mathbb{R}^{m-k}$ and $c'(0)$ not in the tangent space of \mathbb{R}^{m-k} lift to (i.e. using the diffeomorphism property away from \mathbb{R}^{m-k} , can be identified with) curves \tilde{c} in $[\mathbb{R}^m; \mathbb{R}^{m-k}]$ with $\tilde{c}(0) \in \text{ff}$, and two such curves c_j with $c_1(0) = c_2(0)$ satisfy $\tilde{c}_1(0) = \tilde{c}_2(0)$ if and only if $c'_1(0) - c'_2(0)$ is tangent to \mathbb{R}^{m-k} . (This says that invariantly ff is the spherical normal bundle of \mathbb{R}^{m-k} in \mathbb{R}^m , i.e. the quotient of its normal bundle minus its zero section by dilations.)

Let $\hat{X} = \mathbb{R}_x \times \mathbb{R}_y^{n-1}$, $S\hat{X} = \hat{X} \times \mathbb{R} \times \mathbb{S}^{n-2}$. In our setting, as we show momentarily, we start with the map

$$(3.2) \quad \Gamma_+ : S\hat{X} \times [0, \infty) \rightarrow [\hat{X} \times \hat{X}; \text{diag}], \quad \Gamma_+(z, \nu, t) = \gamma_{z, \nu}(t),$$

being a diffeomorphism near $S\hat{X} \times \{0\}$. More precisely, Γ_+ is defined on $\tilde{U} \times (-\tilde{\delta}_0, \tilde{\delta}_0) \times \mathbb{S}^{n-2} \times (-\tilde{\delta}_0, \tilde{\delta}_0)$, and this map is a diffeomorphism onto its range when restricted to a neighborhood of $S\hat{X} \times \{0\}$. To see this, note that the diagonal is the submanifold $z - z' = 0$ of $\hat{X} \times \hat{X}$, so nearby one can use coordinates $z - z' \in \mathbb{R}^n$ (the analogue of w' above) and $z \in \mathbb{R}^n$ (the analogue of w'' above). Thus, coordinates on $[\hat{X} \times \hat{X}; \text{diag}]$ are given by z , $|z' - z|$ and $\frac{z' - z}{|z' - z|}$, and a simple calculation shows that at $t = 0$, one has $\frac{z' - z}{|z' - z|} = \frac{\nu}{|\nu|}$ (with the norms being just Euclidean norms), which proves that Γ_+ as in (3.2) is a diffeomorphism near $S\hat{X} \times \{0\}$. Similarly,

$$(3.3) \quad \Gamma_- : S\hat{X} \times (-\infty, 0] \rightarrow [\hat{X} \times \hat{X}; \text{diag}], \quad \Gamma_-(z, \nu, t) = \gamma_{z, \nu}(t),$$

is a diffeomorphism near $S\hat{X} \times \{0\}$.

Remark 3.1. The analogous results would work with \tilde{X} in place of \hat{X} . Then $S\tilde{X}$ is the sphere bundle of \tilde{X} , i.e. $T\tilde{X} \setminus o$ quotiented out by the \mathbb{R}^+ -action. If we have a Riemannian metric we could take this to be the unit sphere bundle with respect

to this metric, but *any* other choice of a transversal to the dilation orbits in the tangent space of \hat{X} works, such as the unit sphere bundle with respect to another metric, or indeed (locally, in the region of interest) the space of tangent vectors of the form $\lambda \partial_x + \omega \partial_y$, where $\omega \in \mathbb{S}^{n-2}$, considered above.

We now reduce $\delta_0 > 0$ if necessary so that Γ_+ is a diffeomorphism on $U_{x,y} \times (-\delta_0, \delta_0)_\lambda \times \mathbb{S}_\omega^{n-2} \times [0, \delta_0)_t$, and analogously for Γ_- ; we assume this from now on. (Note that in λ_0 we could allow an arbitrary interval with compact closure for this particular purpose.)

Our inversion problem is now that assuming $(If)(x, y, \lambda, \omega) = \int_{\mathbb{R}} f(\gamma_{x,y,\lambda,\omega}(t)) dt$ is known, we would like to recover f from it. (Recall our convention from above; the integral is really over $(-\delta_0, \delta_0)$, and $f(x', y')$ vanishes for $x' \geq x_0/2$.) It is occasionally convenient to assume

$$(3.4) \quad \gamma_{x,y,-\lambda,-\omega}(-t) = \gamma_{x,y,\lambda,\omega}(t).$$

Without this symmetry assumption, we would have two curves with a given tangent line at (x, y) , so having the integral of functions along both, we would have additional information. (In other words, we could simply drop one of these families to arrive at the present setting.)

The idea is simply to average over the family, i.e. to consider for $x > 0$

$$(3.5) \quad (Af)(x, y) = \int_{\mathbb{R}} \int_{\mathbb{S}^{n-2}} (If)(x, y, \lambda, \omega) \tilde{\chi}(x, \lambda) d\lambda d\omega,$$

where $\tilde{\chi}$ is supported in $|\lambda| \leq \sqrt{2C}\sqrt{x}$. One concrete choice that achieves this

$$\tilde{\chi}(x, \lambda) = x^{-1/2} \chi(\lambda/\sqrt{x}),$$

with χ having sufficiently small support near 0; another one is

$$\tilde{\chi}(x, \lambda) = x^{-1} \chi(\lambda/x),$$

where now any compactly supported χ works (for sufficiently small x). We remark that we can allow χ to depend smoothly on ω and y ; over compact sets such a behavior is necessarily uniform since there are no boundaries in these variables.

Remark 3.2. Here we need to recall that γ and A are only locally defined, on some open set O (i.e. γ is defined for $z = (x, y) \in O$ only, and only as long as its image remains in O). However, as we are only interested in applying A to distributions supported in O , and as the ellipticity statements we show are local in nature, this is not a problem. For instance, for K a fixed subset of O , one may replace A by $\psi A \psi$ where $\psi \in C_c^\infty(O)$ is $\equiv 1$ on a neighborhood of K , which is now globally well-defined, and ellipticity statements are unaffected near K .

For any r , we can write A as

$$(3.6) \quad A = L \circ I, \quad L = M_2 \circ \Pi \circ M_1 \circ I,$$

where

$$\Pi u(x, y) = \int_{\mathbb{R}} \int_{\mathbb{S}^{n-2}} u(x, y, \lambda, \omega) d\lambda d\omega,$$

and

$$M_1 u(x, y, \lambda, \omega) = x^r \chi(\lambda/x) u(x, y, \lambda, \omega), \quad (M_2 f)(x, y) = x^{-1-r} f(x, y).$$

Thus, Π is a push-forward map, and thus is bounded on

$$H^s([0, \infty) \times \mathbb{R}^{n-1} \times \mathbb{R} \times \mathbb{S}^{n-2}) \rightarrow H^s([0, \infty) \times \mathbb{R}^{n-1})$$

for all $s \geq 0$, i.e. ‘on H^s ’ in brief, since such a map is bounded on H^s in the absence of boundaries, and there are continuous extension maps from H^s of a half space to H^s of the whole space. On the other hand, as χ is bounded, M_1 is bounded on $L^2(\{x \geq 0\})$ while its j th derivative is bounded by x^{-j} times a constant, so $x^s M_1$ is bounded as map on $H^s(\{x \geq 0\})$ when $s \geq 0$ integer. Thus,

$$L : H^s([0, \infty) \times \mathbb{R}^{n-1} \times \mathbb{R} \times \mathbb{S}^{n-2}) \rightarrow x^{-s-1} H^s([0, \infty) \times \mathbb{R}^{n-1})$$

is bounded. Further, the X-ray transform, I , is itself of the form $I = \tilde{\Pi} \circ \Phi^*$, where Φ^* is pull-back by the map $(z, \nu, t) \mapsto \gamma_{z, \nu}(t)$, $\nu = (\lambda, \omega)$, which has surjective differential in view of the diffeomorphism property of Γ_{\pm} (on the relevant set; recall also that we are assuming that the functions we are applying I to are supported in U), and $\tilde{\Pi}$ is the push-forward given by integration in t . Thus, I itself is bounded

$$I : H^s([0, \infty) \times \mathbb{R}^{n-1}) \rightarrow H^s([0, \infty) \times \mathbb{R}^{n-1} \times \mathbb{R} \times \mathbb{S}^{n-2}).$$

Correspondingly, if we show A is invertible as a map between appropriate spaces of functions supported near $x = 0$ (as discussed in the previous section), concretely weighted Sobolev spaces, with domain space including $H^s([0, \infty) \times \mathbb{R}^{n-1})$ and range space including $x^{-s-1} H^s([0, \infty) \times \mathbb{R}^{n-1})$, we obtain an estimate for f in terms of $I f$ when f satisfies such a support condition and lies in $H^s([0, \infty) \times \mathbb{R}^{n-1})$.

Note that the A defined by (3.5) is certainly a pseudodifferential operator in $x > 0$; moreover, its principal symbol is elliptic if $\chi \geq 0$ with $\chi > 0$ near 0 (this uses $n > 2$) – while this is well-known, we check it below explicitly in the proof of the next proposition. Our main task is to understand the uniform behavior of A to $x = 0$. It turns out that while A itself is not a scattering pseudodifferential operator, its conjugates by exponential weights are:

Proposition 3.3. *Suppose $\chi \in C_c^\infty(\mathbb{R})$. Let $\tilde{\chi}(x, \lambda) = x^{-1} \chi(\lambda/x)$. The operator $A_F = x^{-1} e^{-F/x} A e^{F/x}$ is in $\Psi_{sc}^{-1,0}$ for $F > 0$.*

The main point here regarding the exponential weights is that the Schwartz kernel of A itself is well-behaved near compact subsets of the front face, i.e. where $X = \frac{x'-x}{x^2}$ and $Y = \frac{y'-y}{x}$ are bounded, but is not so well-behaved as $(X, Y) \rightarrow \infty$. However, the support conditions on χ insure that $X \rightarrow +\infty$ on the support of the Schwartz kernel of A (with a suitable estimate), and thus the exponential conjugation gives exponential decay of the *conjugated kernel* as $(X, Y) \rightarrow \infty$, giving the conclusion of the proposition.

Proof. At first work in $x > 0$, ignoring the limit $x \rightarrow 0$. Then it is standard that A is a pseudodifferential operator (the weights are harmless then), but it is instructive to prove this in a manner that extends seamlessly to the general case.

With Γ_{\pm} as in (3.2)-(3.3), for $\tilde{\chi}$ an arbitrary smooth function on $S\tilde{X}$ (not necessarily dependent just on x, λ) the diffeomorphism property on $S\tilde{X} \times [0, \delta_0)$ allows one to rewrite, with $|d\nu|$ denoting a smooth measure on the transversal such as $|d\lambda| |d\omega|$,

$$A f(z) = \sum_{\bullet=+,-} \int f(z') \tilde{\chi}(\Gamma_{\bullet}^{-1}(z, z')) (\Gamma_{\bullet}^{-1})^* (|d\nu| dt)$$

in terms of z, z' as

$$(3.7) \quad \begin{aligned} & \int f(z')|z' - z|^{-n+1} b\left(z, \frac{z' - z}{|z' - z|}, |z' - z|\right) dz', \\ & b\left(z, \frac{z' - z}{|z' - z|}, 0\right) = \tilde{\chi}\left(z, \frac{z' - z}{|z' - z|}\right) \sigma\left(z, \frac{z' - z}{|z' - z|}\right), \end{aligned}$$

where $\sigma > 0$ is bounded below – it is the change of variables Jacobian factor. The two terms Γ_{\pm} are in fact identical by the symmetry assumption on γ , (3.4), so we can ignore Γ_- . In particular, A is a pseudodifferential operator with principal symbol given by the Fourier transform of

$$|z' - z|^{-n+1} b\left(z, \frac{z' - z}{|z' - z|}, 0\right) = |z' - z|^{-n+1} (\tilde{\chi}\sigma)\left(z, \frac{z' - z}{|z' - z|}\right)$$

in $Z = z' - z$. One can insert a cutoff ϕ in $|Z|$ with compact support, identically 1 near 0 (considered as an even function on \mathbb{R}), without changing the result modulo rapid decay, i.e. as a principal symbol, the result is not affected. The latter can be computed easily as

$$\begin{aligned} & \int_{\mathbb{R}^n} e^{-iZ \cdot \zeta} |Z|^{-n+1} (\tilde{\chi}\sigma)(z, \hat{Z}) \phi(|Z|) dZ = \int_0^{\infty} \int_{\mathbb{S}^{n-1}} e^{-it\hat{Z} \cdot \zeta} (\tilde{\chi}\sigma)(z, \hat{Z}) \phi(t) dt d\hat{Z} \\ & = \frac{1}{2} \int_{\mathbb{R}} \int_{\mathbb{S}^{n-1}} e^{-it\hat{Z} \cdot \zeta} (\tilde{\chi}\sigma)(z, \hat{Z}) \phi(t) dt d\hat{Z} = \frac{1}{2} \int_{\mathbb{S}^{n-1}} \hat{\phi}(\hat{Z} \cdot \zeta) (\tilde{\chi}\sigma)(z, \hat{Z}) d\hat{Z}; \end{aligned}$$

here $\hat{\phi}$ is the Fourier transform of ϕ . Fixing $\hat{\zeta} = \frac{\zeta}{|\zeta|}$, since $\hat{\phi}$ is Schwartz, if $\tilde{\chi}\sigma$ is supported away from the equatorial sphere $\{\hat{Z} : \hat{Z} \cdot \hat{\zeta} = 0\}$, $|\hat{Z} \cdot \hat{\zeta}| > c|\zeta|$ on its support for some $c > 0$, and then for all N , $\hat{\phi}(\hat{Z} \cdot \zeta) \leq \tilde{C}_N |\zeta|^{-N}$, and thus we conclude that the integral is Schwartz and thus gives no contribution to the principal symbol. Correspondingly (by using a partition of unity), it suffices to consider a neighborhood of the equator and assume $\tilde{\chi}\sigma$ is supported here. Then one can write $Z = (Z^{\parallel}, Z^{\perp})$ according to the orthogonal decomposition relative to $\hat{\zeta} = \frac{\zeta}{|\zeta|}$, so $Z^{\parallel} = Z \cdot \hat{\zeta}$, similarly for \hat{Z} , and $d\hat{Z}$ is of the form $a(\hat{Z}^{\parallel}) d\hat{Z}^{\parallel} d\theta$, $\theta = \frac{\hat{Z}^{\perp}}{|\hat{Z}^{\perp}|} \in \mathbb{S}^{n-2}$ with $a(0) = 1$ since $\hat{Z}^{\perp} = (1 - |\hat{Z}^{\parallel}|^2)^{1/2} \theta$. Thus, one has

$$\begin{aligned} & \frac{1}{2} \int_{\mathbb{R}} \int_{\mathbb{S}^{n-2}} \hat{\phi}(\hat{Z}^{\parallel} |\zeta|) (\tilde{\chi}\sigma)(z, \hat{Z}^{\parallel} \hat{\zeta} + (1 - |\hat{Z}^{\parallel}|^2)^{1/2} \theta) a(\hat{Z}^{\parallel}) d\theta d\hat{Z}^{\parallel} \\ & = \frac{1}{2|\zeta|} \int_{\mathbb{R}} (|\zeta| \hat{\phi}(\hat{Z}^{\parallel} |\zeta|)) a(\hat{Z}^{\parallel}) \left(\int_{\mathbb{S}^{n-2}} (\tilde{\chi}\sigma)(z, \hat{Z}^{\parallel} \hat{\zeta} + (1 - |\hat{Z}^{\parallel}|^2)^{1/2} \theta) d\theta \right) d\hat{Z}^{\parallel} \end{aligned}$$

Since $(|\zeta| \hat{\phi}(\hat{Z}^{\parallel} |\zeta|)) \rightarrow \delta_0$ in distributions as $|\zeta| \rightarrow \infty$, this is $|\zeta|^{-1} \int_{\mathbb{S}^{n-2}} (\tilde{\chi}\sigma)(z, \theta) d\theta$ modulo terms decaying faster as $|\zeta| \rightarrow \infty$; indeed, one easily sees by expanding $\tilde{\chi}\sigma$ around $\hat{Z}^{\parallel} = 0$ that this asymptotic holds modulo $O(|\zeta|^{-2})$ terms. In other words, the principal symbol of A at (z, ζ) is a constant multiple of

$$(3.8) \quad |\zeta|^{-1} \int (\tilde{\chi}\sigma)(z, \hat{Z}^{\perp}) dZ^{\perp}.$$

In particular, if $\tilde{\chi} \geq 0$, then as long as for each (z, ζ) , $\zeta \neq 0$, there is \hat{Z} perpendicular to ζ with $\tilde{\chi}$ non-zero at (z, \hat{Z}) , then A is an elliptic order -1 pseudodifferential operator, in accordance with the results of Stefanov and Uhlmann [16]. This is indeed the case with our choice of $\tilde{\chi}$, provided $n > 2$.

We now turn to the scattering behavior, i.e. as at least one of $x, x' \rightarrow 0$. Note that from (3.1), on the support of $\tilde{\chi}$, $x' \geq x - c_0 x^2$, for x small. We in fact show below that on the support of $\tilde{\chi}$, X is bounded below, and $X \rightarrow +\infty$ if $|Y| \rightarrow \infty$, and indeed $X \geq C_1 |Y|^2$ for $|Y|$ sufficiently large, $C_1 > 0$. Here we recall from Section 2 that

$$X = \frac{x - x'}{x^2}, \quad Y = \frac{y - y'}{x}.$$

With K denoting the Schwartz kernel of A , as

$$x^{-1} - (x')^{-1} = \frac{x' - x}{xx'} = X \frac{x}{x'} = X/(1 + xX),$$

A_F has Schwartz kernel

$$(3.9) \quad \begin{aligned} K^b(x, y, X, Y) &= x^{-1} e^{-F(x^{-1} - (x')^{-1})} K(x, y, X, Y) \\ &= x^{-1} e^{-F X/(1+xX)} K(x, y, X, Y). \end{aligned}$$

Taking into account the polynomial bounds on K in terms of X, Y , and $x' \geq x - c_0 x^2$ implying that X is bounded below as shown later in the proof, further that $X \rightarrow +\infty$ as $|Y| \rightarrow \infty$ with $X \geq C_1 |Y|^2$, exponential decay of K^b as well as its derivatives follows easily for $F > 0$. Thus, the main claim is that K^b is smooth for (X, Y) finite, non-zero, conormal to $(X, Y) = 0$.

Now, on $\Gamma_+(\text{supp } \tilde{\chi} \times [0, \delta_0])$, $|x - x'| \leq C|y - y'|$ means that locally in this region $x, y, |y' - y|, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|}$ are coordinates on $[\tilde{X} \times \tilde{X}; \text{diag}]$ – indeed, this corresponds to using the transversal $|y' - y| = 1$ to dilations in $\mathbb{R}^n = \mathbb{R}_{x' - x} \times \mathbb{R}_{y' - y}^{n-1}$ in place of the unit sphere $|(x' - x, y' - y)| = 1$, which is indeed a transversal where $y' - y$ is large relative to $x' - x$, i.e. in our region of interest. Further, $\Gamma_+(x, y, \lambda, \omega, 0)$ is, in terms of these coordinates, $(x, y, \lambda, \omega, 0)$. In general, thus,

$$\lambda \left(\Gamma_+^{-1} \left(x, y, |y' - y|, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|} \right) \right) = \Lambda \left(x, y, |y' - y|, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|} \right),$$

with

$$\Lambda \left(x, y, 0, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|} \right) = \frac{x' - x}{|y' - y|},$$

so (suppressing Γ_+ on the left hand side)

$$\lambda = \frac{x' - x}{|y' - y|} + |y' - y| \tilde{\Lambda} \left(x, y, |y' - y|, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|} \right).$$

Now, in terms of the scattering coordinates,

$$|y' - y| = x|Y|, \quad \frac{x' - x}{|y' - y|} = \frac{xX}{|Y|}, \quad \frac{y' - y}{|y' - y|} = \hat{Y},$$

so (suppressing Γ_+ composed with the scattering blow up map on the left hand side)

$$(3.10) \quad \frac{\lambda}{x} = \frac{X}{|Y|} + |Y| \tilde{\Lambda} \left(x, y, x|Y|, \frac{xX}{|Y|}, \hat{Y} \right).$$

Similarly,

$$\begin{aligned} \omega &= \frac{y' - y}{|y' - y|} + |y' - y| \tilde{\Omega} \left(x, y, |y' - y|, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|} \right) \\ &= \hat{Y} + x|Y| \tilde{\Omega} \left(x, y, x|Y|, \frac{xX}{|Y|}, \hat{Y} \right) \end{aligned}$$

and

$$\begin{aligned}
 (3.11) \quad t &= |y' - y| + |y' - y|^2 \tilde{T}\left(x, y, |y' - y|, \frac{x' - x}{|y' - y|}, \frac{y' - y}{|y' - y|}\right) \\
 &= x|Y| + x^2|Y|^2 \tilde{T}\left(x, y, x|Y|, \frac{xX}{|Y|}, \hat{Y}\right).
 \end{aligned}$$

Thus,

$$\begin{aligned}
 (3.12) \quad dt \, d\lambda \, d\omega &= J\left(x, y, \frac{X}{|Y|}, |Y|, \hat{Y}\right) x^2 |Y|^{-1} dX \, d|Y| \, d\hat{Y} \\
 &= J\left(x, y, \frac{X}{|Y|}, |Y|, \hat{Y}\right) x^2 |Y|^{-n+1} dX \, dY
 \end{aligned}$$

where the density factor J is smooth and positive, and $J|_{x=0} = 1$. Also, on the blow-up of the scattering diagonal, $\{X = 0, Y = 0\}$, in the region $|Y| > \epsilon|X|$, thus on the support of χ in view of (3.10),

$$x, y, |Y|, \frac{X}{|Y|}, \hat{Y}$$

are valid coordinates, with $|Y|$ being the defining function of the front face of this blow up (i.e. of the lifted diagonal). Taking into account the x^{-1} in the definition in $\tilde{\chi}$, we thus deduce that K^b is given by

$$(3.13) \quad e^{-FX/(1+X)} \chi\left(\frac{X}{|Y|} + |Y| \tilde{\Lambda}(x, y, x|Y|, \frac{xX}{|Y|}, \hat{Y})\right) |Y|^{-n+1} J\left(x, y, \frac{X}{|Y|}, |Y|, \hat{Y}\right),$$

so in particular it is conormal to the front face on the blow-up of the scattering diagonal, of the form $\rho^{-n+1}b$, where b is smooth up to the front face, and without the first exponential factor it, together with its derivatives (in x, y, X, Y) has polynomial growth estimates as $(X, Y) \rightarrow \infty$, i.e. the derivatives satisfy bounds $\leq C|(X, Y)|^N$ for some C, N (depending on the derivative). Decomposing K^b into pieces supported in, say, $|(X, Y)| < 2$ and $|(X, Y)| > 1$ by a partition of unity, we show in the next paragraph that the latter is Schwartz in (X, Y) due the exponential decay of the first factor of (3.13) on the support of the second factor. On the other hand, for the former term, supported in $|(X, Y)| < 2$, calculations as in (3.7) in Fourier transforming this in (X, Y) show that this term of K^b is indeed the Schwartz kernel of an element of $\Psi_{\text{sc}}^{-1,0}$, with standard principal symbol being given by the analogue of (3.8). Here the additional information is in the behavior at $x = 0$, but given that our operator *is* an element of $\Psi_{\text{sc}}^{-1,0}$, the same information can be obtained from computing the boundary principal symbol, which we need in any case.

We use (3.11) to express $\frac{\lambda}{x}$ using

$$(3.14) \quad x' = x + \lambda t + \alpha(x, y, \lambda, \omega)t^2 + O(t^3), \quad y' = y + \omega t + O(t^2),$$

where the $O(t^2)$ and $O(t^3)$ terms have coefficients which are smooth in (x, y, λ, ω) . Thus,

$$X = \frac{x' - x}{x^2} = \frac{\lambda t}{x^2} + \frac{\alpha t^2}{x^2} + \frac{t^3}{x^2} \Upsilon(x, y, x\mu, \omega, t),$$

with Υ a smooth function of its arguments, so

$$X = \frac{\lambda(\Gamma_+^{-1})}{x} |Y| (1 + x|Y| \tilde{T}(x, y, x|Y|, xX/|Y|, \hat{Y})) \\ + \alpha(\Gamma_+^{-1}) |Y|^2 (1 + x|Y| \tilde{T}(x, y, x|Y|, xX/|Y|, \hat{Y}))^2 + x|Y|^3 \Upsilon(\Gamma_+^{-1}),$$

and thus

$$(3.15) \quad \frac{\lambda(\Gamma_+^{-1})}{x} = \frac{X - \alpha(\Gamma_+^{-1})|Y|^2}{|Y|} + O(x),$$

where the $O(x)$ has smooth coefficients in terms of $x, y, x|Y|, xX/|Y|, \hat{Y}$. Thus, for $\mu = \frac{\lambda}{x} \in (-c, c)$, $-c|Y| < X - \alpha(\Gamma_+^{-1})|Y|^2 < c|Y|$, which shows (by the positive definiteness of α) that $X \rightarrow +\infty$ on $\text{supp } \tilde{\chi}$ if $|Y| \rightarrow \infty$, and indeed, for $|Y|$ sufficiently large, $X > C_1|Y|^2$ for some $C_1 > 0$.

As already explained, this proves the proposition, since now for all N' the exponential factor in (3.13) is $\leq C'|(X, Y)|^{-N'}$ for suitable C' on the support of the second factor, so combined with the polynomial estimates for the derivatives of the second and third factors, it follows that K^b is smooth in (x, y) , with values in functions Schwartz in (X, Y) for $(X, Y) \neq 0$, and conormal to $(X, Y) = 0$, which is exactly the characterization of the Schwartz kernel of a scattering pseudodifferential operator. \square

Remark 3.4. We now explain the form these arguments would take for a different scaling chosen for $\tilde{\chi}$. By (3.14) for $\lambda = \sqrt{x}\mu$, with μ in a compact set near 0 (i.e. the first, $O(\sqrt{x})$ localization used above), $x' = x + \sqrt{x}\mu t + \alpha(x, y, \sqrt{x}\mu, \omega, t)t^2 + O(t^3)$ gives

$$x' \leq C'(x + |y - y'|^2),$$

indicating that $|y' - y|/x^{1/2}$ is the appropriate homogeneous variable for analysis; using $X = \frac{\sqrt{x'} - \sqrt{x}}{\sqrt{x}}$, $Y = \frac{y' - y}{\sqrt{x}}$, this amounts to a statement that the analysis is well-behaved on the 0 double space of Mazzeo-Melrose [9] when the smooth structure is given by the boundary defining function \sqrt{x} . This is a somewhat complicated space with a non-commutative normal operator at infinity; there's a reduced normal operator after a partial Fourier transform and rescaling which is a b-scattering (or Bessel) type pseudodifferential operator on a half-line. This is the reason for using our sharper cutoff, which puts us into the more amenable setting of Melrose's scattering calculus, as described above.

We now compute the boundary principal symbol of A_F . Indeed, this is immediate from (3.13) and (3.15) which show that at $x = 0$ (i.e. the scattering front face) the Schwartz kernel of A_F is

$$e^{-FX} |Y|^{-n+1} \chi\left((X - \alpha(0, y, 0, \hat{Y})|Y|^2)/|Y|\right) = \tilde{K}(y, X, Y).$$

As described in Section 2, for each y , $\tilde{K}(y, \cdot, \cdot)$ acts as a convolution operator, thus it becomes a multiplication operator upon Fourier transforming in (X, Y) , and the desired invertibility $H_{sc}^{s,r} \rightarrow H_{sc}^{s+1,r}$ amounts to the Fourier transformed kernel, $\hat{K}(y, \cdot, \cdot)$ being bounded below in absolute value by $c\langle(\xi, \eta)\rangle^{-1}$, $c > 0$ (here (ξ, η) are the Fourier dual variables of (X, Y)). Thus, we need to compute the inverse

Fourier transform of $\tilde{K}(y, \cdot, \cdot)$, and find F such that the desired bound holds. Note also that if χ depends on y and ω as discussed above, we simply have

$$\chi\left((X - \alpha(0, y, 0, \hat{Y})|Y|^2)/|Y|, y, \hat{Y}\right)$$

in the above expression for the Schwartz kernel at the front face. We have thus shown

Lemma 3.5. *The boundary principal symbol of $x^{-1}e^{-F/x}Ae^{F/x}$ is the (X, Y) -Fourier transform of*

$$\tilde{K}(y, X, Y) = e^{-FX}|Y|^{-n+1}\chi\left((X - \alpha(0, y, 0, \hat{Y})|Y|^2)/|Y|, y, \hat{Y}\right).$$

In order to find a suitable χ , we first make a slightly inadmissible choice for an easier computation, namely we take $\chi(s) = e^{-s^2/(2\nu)}$ with ν to be fixed (and allowed to depend on y and \hat{Y}), so $\hat{\chi}(\cdot) = c\sqrt{\nu}e^{-\nu|\cdot|^2/2}$ for appropriate $c > 0$. Thus, χ does not have compact support, and an approximation argument will be necessary. Now, in general (for arbitrary χ which has superexponential decay so its Fourier transform is entire), the Fourier transform in X is

$$(3.16) \quad \mathcal{F}_X \tilde{K}(y, \xi, Y) = |Y|^{2-n} e^{-\alpha F|Y|^2} e^{-i\alpha\xi|Y|^2} \hat{\chi}((\xi - iF)|Y|),$$

as follows by taking into account the effect of translations, dilations and multiplication by exponential weights on the Fourier transform (the last two of which are closely related). Here α is a function of y and \hat{Y} , as above. Substituting the particular χ yields a non-zero multiple of

$$(3.17) \quad \begin{aligned} & \sqrt{\nu}|Y|^{2-n} e^{-\alpha F|Y|^2} e^{-i\alpha\xi|Y|^2} e^{-\nu(\xi - iF)^2|Y|^2/2} \\ & = \sqrt{\nu}|Y|^{2-n} e^{-(2\alpha F + \nu\xi^2 - \nu F^2)|Y|^2/2} e^{-i(\alpha - F\nu)\xi|Y|^2}. \end{aligned}$$

Now, the Y -Fourier transform of $|Y|^{2-n}$ is a homogeneous radial (i.e. $\text{SO}(n-1)$ -invariant) function of order -1 , so it is a non-zero multiple of $|\eta|^{-1}$, with η the Fourier-dual variable of Y . Notice that this uses very strongly that we have $n > 2$; $n = 2$ would give a delta distribution. Thus, if the Y -Fourier transform of

$$e^{-(2\alpha F + \nu\xi^2 - \nu F^2)|Y|^2/2} e^{-i(\alpha - F\nu)|Y|^2}$$

is positive, then the Fourier transform of the product, which is given by the convolution (in η) of these, is also positive, and with asymptotic behavior given by that of $|\eta|^{-1}$ provided the Fourier transform of the Gaussian is Schwartz. Indeed, if one Fourier transforms $|Y|^{2-n}\psi(y, Y)$, where ψ is Schwartz in the last variable, only the behavior near $Y = 0$ contributes to the asymptotics as $\eta \rightarrow \infty$, and thus using the Taylor series of ψ , one obtains the asymptotic expansion of the Fourier transform as $\eta \rightarrow \infty$ as a classical polyhomogeneous function (with the expansion in terms of $|\eta|^{-1-j}$, $j \geq 0$ integer).

So assume now that α is a positive definite quadratic form in \hat{Y} and take $\nu = F^{-1}\alpha$ (so same holds for ν , i.e. ν is a quadratic form in \hat{Y}). Thus, one has $\alpha|Y|^2 = Q(Y, Y)$, a quadratic form in Y . Thus, writing $Q^{-1}(Y, Y)$ for the dual quadratic form, and taking $\chi(s) = e^{-s^2/(2F^{-1}Q(\hat{Y}, \hat{Y}))}$, we have

$$\hat{\chi}(\sigma) = c(F^{-1}Q(\hat{Y}, \hat{Y}))^{1/2} e^{-F^{-1}Q(\sigma\hat{Y}, \sigma\hat{Y})/2}.$$

In view of (3.16)-(3.17), $\mathcal{F}_X \tilde{K}(y, \xi, Y)$ is a non-zero multiple of

$$\begin{aligned} & \sqrt{\nu} |Y|^{2-n} e^{-\alpha F |Y|^2} e^{-i\alpha \xi |Y|^2} e^{-\nu(\xi-iF)^2 |Y|^2/2} \\ &= \sqrt{\nu} |Y|^{2-n} e^{-(2\alpha F + \nu \xi^2 - \nu F^2) |Y|^2/2} e^{-i(\alpha - F\nu)\xi |Y|^2} \\ &= F^{-1/2} \sqrt{\alpha} |Y|^{2-n} e^{-(\xi^2 + F^2)\alpha F^{-1} |Y|^2/2} = F^{-1/2} \sqrt{\alpha} |Y|^{2-n} e^{-F^{-1}(\xi^2 + F^2)Q(Y,Y)/2}, \end{aligned}$$

where the last factor is a real Gaussian since the oscillatory factor in (3.17) becomes identically 1. This is Schwartz in Y for $F > 0$, and thus the Fourier transform is a positive multiple of

$$(\det Q)^{-1/2} F^{(n-1)/2} (\xi^2 + F^2)^{-(n-1)/2} e^{-F Q^{-1}(\eta, \eta)/(2(\xi^2 + F^2))}.$$

which satisfies the requirements from the previous paragraph (positive Schwartz function).

One has to be a bit careful about the joint (ξ, η) -behavior, i.e. when ξ is also going to infinity, and where we still need lower bounds. The Fourier transform of $|Y|^{2-n} e^{-F^{-1}(\xi^2)Q(Y,Y)/2}$, with $\langle \xi \rangle = (\xi^2 + F^2)^{1/2}$, is a constant multiple of

$$\begin{aligned} & \int |\eta - \zeta|^{-1} \langle \xi \rangle^{-(n-1)} e^{-F Q^{-1}(\zeta, \zeta)/(2\langle \xi \rangle^2)} d\zeta \\ &= \int |\eta - \langle \xi \rangle \zeta'|^{-1} e^{-F Q^{-1}(\zeta', \zeta')/2} d\zeta' \\ &= \langle \xi \rangle^{-1} \int |\eta/\langle \xi \rangle - \zeta'|^{-1} e^{-F Q^{-1}(\zeta', \zeta')/2} d\zeta' = \langle \xi \rangle^{-1} \varphi(\eta/\langle \xi \rangle) \end{aligned}$$

where we wrote $\zeta' = \zeta/\langle \xi \rangle$, and where φ is an elliptic positive classical symbol of order -1 , namely the convolution of $|\cdot|^{-1}$ with the Schwartz function $e^{-F Q^{-1}(\cdot, \cdot)/2}$. This assures lower bounds $c\langle(\xi, \eta)\rangle^{-1}$, $c > 0$, i.e. elliptic lower bounds. Indeed, this is immediate when $\langle \xi \rangle > \epsilon|\eta|$, for $\langle \xi \rangle^{-1}$ is equivalent to $\langle(\xi, \eta)\rangle^{-1}$ in this region in terms of decay rates, while $\varphi(\eta/\langle \xi \rangle)$ is a 0th order symbol in this region. To see what happens when $|\eta| > \epsilon\langle \xi \rangle$, notice that by virtue of the classicality in fact have $\varphi(\eta') = |\eta'|^{-1} \tilde{\varphi}(\langle \xi \rangle/|\eta'|, \eta'/|\eta'|)$, with $\tilde{\varphi}$ smooth near 0 in the first argument. Thus, we obtain

$$\langle \xi \rangle^{-1} \varphi(\eta/\langle \xi \rangle) = |\eta|^{-1} \tilde{\varphi}(\langle \xi \rangle/|\eta|, \eta'/|\eta'|),$$

which is a symbol of order -1 in $|\xi| \leq C|\eta|$, and $|\eta|^{-1}$ is equivalent to $\langle(\xi, \eta)\rangle^{-1}$ here, completing the proof of the ellipticity claim.

Now, if χ is not a Gaussian, but rather one has a sequence χ_n in $C_c^\infty(\mathbb{R})$ which converges to the Gaussian in Schwartz functions (notice that this does not imply that the Fourier transform of χ_n is pointwise positive for any n , which is the reason we need to use the Fourier transform of the Gaussian directly), then the Fourier transforms converge in the appropriate spaces, which suffices to conclude that for sufficiently large n , letting $\chi = \chi_n$, the Fourier transform \hat{K} , i.e. the boundary principal symbol, still has lower bounds $\tilde{C}\langle(\xi, \eta)\rangle^{-1}$, $\tilde{C} > 0$, as desired. We have thus proved:

Lemma 3.6. *For $F > 0$ there exists $\chi \in C_c^\infty(\mathbb{R})$, $\chi \geq 0$, $\chi(0) = 1$, such that for the corresponding operator $x^{-1}e^{-F/x}Ae^{F/x}$ the boundary symbol is elliptic; indeed, this holds for all χ sufficiently close in Schwartz space to a specific Gaussian.*

Hence, we have

$$B = x^{-1}e^{-F/x}Ae^{F/x} \in \Psi_{\text{sc}}^{-1,0}$$

elliptic both in the sense of the standard principal symbol (in the set of interest O), and the scattering principal symbol, which is at $x = 0$, and in particular the results of Section 2 are applicable. Thus, elements of the kernel of A which have support in the compact subset K of O is finite dimensional, and further a stability estimate holds on a complementary subspace of this finite dimensional subspace. Further, with $x = x_c = \tilde{x} + c$, as discussed at the beginning of this section, the arguments at the end of Section 2 show that for sufficiently small c , this subspace of the kernel of $A = A_c$ is actually trivial, and one has a stability estimate in $M_c = \{x_c > 0\}$ for functions supported in K . Thus, for c small, writing the support condition as final subscript,

$$A = xe^{F/x} B e^{-F/x} : e^{F/x} H_{\text{sc}}^{s,r}(M_c)_K \rightarrow xe^{F/x} H_{\text{sc}}^{s+1,r}(M_c) = e^{F/x} H_{\text{sc}}^{s+1,r+1}(M_c)$$

satisfies estimates

$$\|f\|_{e^{F/x} H_{\text{sc}}^{s,r}(M_c)_K} \leq C \|Af\|_{e^{F/x} H_{\text{sc}}^{s+1,r+1}(M_c)}.$$

In particular, if one is willing to give up polynomial weights as unimportant at the cost of losing δ/x in the exponential weight, $\delta > 0$, and one uses that for $s \geq 0$, $H^s(M_c) \subset H_{\text{sc}}^{s,r}(M_c)$ for $r \leq -\frac{n+1}{2}$ while for $r \geq -\frac{n+1}{2} + 2s$, $H_{\text{sc}}^{s,r}(M_c) \subset H^s(M_c)$, with continuous inclusion maps, see (2.2)-(2.3), we have

$$\|f\|_{e^{(F+\delta)/x} H^s(M_c)_K} \leq C \|Af\|_{e^{F/x} H^{s+1}(M_c)}.$$

Now, using the decomposition $A = L \circ I$ of A in (3.6), and the boundedness statements following it, we have for all $F > 0$,

$$\|Af\|_{e^{F/x} H^{s+1}(M_c)} \leq C' \|If\|_{H^{s+1}(\mathcal{M}_{M_c})}$$

when $f \in H^{s+1}(M_c)_K$. For the convenience of the reader, we summarize all the maps for $F > 0$ and with $r \in \mathbb{R}$, $\delta > 0$ arbitrary, in a commutative diagram:

$$\begin{array}{ccccc} e^{F/x} H_{\text{sc}}^{s,r}(M_c)_K & \xrightarrow{A} & e^{F/x} H_{\text{sc}}^{s+1,r+1}(M_c) & \xrightarrow{G} & e^{F/x} H_{\text{sc}}^{s,r}(M_c) \\ \uparrow & & \uparrow & & \downarrow \\ H^{s+1}(M_c)_K & \xrightarrow{L \circ I} & H^{s+1}(M_c) & & e^{(F+\delta)/x} H^s(M_c) \end{array}$$

with all vertical arrows inclusion maps, G the inverse of A on the range of A , and with

$$H^{s+1}(M_c)_K \xrightarrow{I} H^{s+1}(\mathcal{M}_{O_c}) \xrightarrow{L} H_{\text{sc}}^{s+1}(M_c)$$

being the lower left composite map $L \circ I$ written out in detail. Indeed, note that even $s = -1$ is allowed with the inclusions we stated; factoring the first inclusion via $H^{s+1}(M_c) \rightarrow e^{F/x} H_{\text{sc}}^{s+1,r}(M_c) \rightarrow e^{F/x} H_{\text{sc}}^{s,r}(M_c)$ proves it since $s+1 \geq 0$, while using (2.4) (with $-s$ in place of s) gives the last inclusion map. In combination, this completes the proof of the main theorem for $s \geq -1$ in the notation here (thus $s \geq 0$ for the notation of the main theorem), with F replaced by $F + \delta$ – as both $F > 0$ and $\delta > 0$ are arbitrary, this means that the original statement is proved.

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